

# Chapter 17 WAVES II

### 1. Sound Waves

Sound waves are longitudinal mechanical waves that can travel through solids, liquids and gases. We focus in this chapter on sound waves that travel through air and that are audible to people.

In the figure, point *S* represents a tiny sound source, or a **point source**, that emits sound in all directions.

Wavefronts are surfaces over which the sound oscillations (or displacements) due to the sound wave have the same value. Wavefronts are represented by whole or partial circles.



### 1. Sound Waves

**Rays** are directed lines perpendicular to the wavefronts that indicate the direction of travel of the wavefronts.

The wavefronts in the figure are spheres, and therefore the waves are said to be **spherical**.

Far from the source, the wavefronts can be approximated as planes, and the waves are said to be **planar**.



# 2. The Speed of Sound

The speed of any mechanical wave depends on both an inertial property of the medium (to store kinetic energy) and an elastic property of the medium (to store potential energy).

We can thus generalize the equation for the speed of a wave along a string by writing

$$\nu = \sqrt{\frac{\tau}{\mu}} = \sqrt{\frac{\text{elastic propery}}{\text{inertial property}}}$$

Because the medium of a sound wave is air, we replace  $\mu$  with  $\rho$ , the density of air.

In a string, the potential energy is associated with the stretching of the string elements. In air, potential energy is associated with compressions and expansions of volume elements of the air.

# 2. The Speed of Sound

The property that determines the extent to which an element of a medium changes in volume when the pressure on it changes is the **bulk modulus** *B*:

$$B = -\frac{\Delta p}{\Delta V/V}.$$

Replacing  $\tau$  with B yields the expression for the speed of sound:

$$v = \sqrt{\frac{B}{\rho}}.$$

#### The Speed of Sound

Medium	Speed (m/s)		
Gases			
Air $(0^{\circ}C)$	331		
Air (20°C)	343		
Helium	965		
Hydrogen	1284		
Liquids			
Water (0°C)	1402		
Water (20°C)	1482		
Seawater	1522		
Solids			
Aluminum	6420		
Steel	5941		
Granite	6000		

The oscillations of air elements due to a traveling sound wave are similar to those of a string's elements due to a transverse wave. However, air elements oscillate longitudinally rather than transversely.

The displacements s(x, t) of air elements are given by

$$s(x,t) = s_m \cos(kx - \omega t),$$

where  $s_m$  is the **displacement amplitude**.



The air pressure varies sinusoidally as the wave moves:

 $\Delta p(x,t) = \Delta p_m \sin(kx - \omega t),$ 

where  $\Delta p_m$  is the **pressure amplitude**; the maximum increase or decrease in pressure. It is given by

 $\Delta p_m = (v\rho\omega)s_m.$ 

Usually,  $\Delta p_m$  is very much less than the pressure where there is no wave.



Note that negative  $\Delta p$  corresponds to an expansion and positive  $\Delta p$  corresponds to compression.

s and  $\Delta p$  are  $\frac{\pi}{2}$  out of phase, as shown in the figure.



**Example 1**: The maximum pressure amplitude  $\Delta p_m$  that the human ear can tolerate in loud sounds is about 28 Pa (which is very much less than the normal air pressure of about 10<sup>5</sup> Pa). What is the displacement amplitude  $s_m$  for such a sound in air of density  $\rho = 1.21 \text{ kg/m}^3$ , at a frequency of 1000 Hz and a speed of 343 m/s?

We know that  $\Delta p_m = (v\rho\omega)s_m$ , which gives

 $s_m = \frac{\Delta p_m}{\nu \rho \omega} = \frac{28 \text{ Pa}}{(343 \text{ m/s})(1.21 \text{ kg/m}^3)[2\pi (1000 \text{ Hz})]} = 1.1 \times 10^{-5} \text{ m} = 11 \text{ \mum}.$ 

Consider the situation shown in the figure. Two point sources  $S_1$  and  $S_2$  emit sound waves that are in phase and of the same wavelength  $\lambda$ . We assume that the distances  $L_1$  and  $L_2$  are much larger than that between the sources. Thus we can approximate the waves at P to be in the same direction.

If  $L_1 = L_2$ , the two waves would be in phase at Pand their interference there would be fully constructive. However,  $L_1$  is different from  $L_2$  in general. The phase difference  $\phi$  at P depends on the **path length difference**  $\Delta L = |L_2 - L_1|$ .

We use the fact that a phase difference of  $2\pi$  corresponds to one wavelength. Thus, we write

$$\frac{\phi}{2\pi} = \frac{\Delta L}{\lambda}$$

or

$$\phi = 2\pi \frac{\Delta L}{\lambda}.$$

Fully constructive interference occurs when

$$\phi = 2\pi m$$
 ,  $m = 0, 1, 2, ...$ 

which corresponds to

$$\frac{\Delta L}{\lambda} = 0, 1, 2, \dots$$

Fully destructive interference occurs when

$$\phi = \pi (2m + 1)$$
  $m = 0, 1, 2, ...$ 

which corresponds to

$$\frac{\Delta L}{\lambda} = \frac{1}{2}, \frac{3}{2}, \frac{5}{2}, \dots \qquad \phi = 2\pi \frac{\Delta L}{\lambda}.$$

The interference is intermediate for other values of  $\Delta L/\lambda$ .

**Example 2**: In the figure, two point sources  $S_1$  and  $S_2$ , which are in phase and separated by distance  $D = 1.5\lambda$ , emit identical sound waves of wavelength  $\lambda$ .

(a) What is the path length difference of the waves from  $S_1$  and  $S_2$  at point  $P_1$ , which lies on D/2 the perpendicular bisector of distance D, at a distance greater than D from the sources? What type of interference occurs at  $P_1$ ?

 $\Delta L=0.$ 

The interference is therefore fully constructive.



(b) What are the path length difference and type of interference at point  $P_2$ ?

 $\Delta L = 1.5 \lambda.$ 

The interference is therefore fully destructive.



(c) The figure shows a circle with a radius much greater than D, centered on the midpoint between sources  $S_1$  and  $S_2$ . What is the number of points N around this circle at which the interference is fully constructive? (That is, at how many points do the waves arrive exactly in phase?)

At point a,  $\Delta L = 0$ . It increases as you go down along the circle until it becomes  $1.5 \lambda$  at point d. Therefore, there is point between a and d at which  $\Delta L = \lambda$ , where a fully constructive interference occurs.



Thus, there are six points (N = 6) at which the interference is fully constructive.



The intensity *I* of a sound wave at a surface is the average rate per unit area at which energy is transferred by the wave through or onto the surface:

$$I=\frac{P}{A},$$

where P is the average power (time rate of energy transfer) of the wave and A is the area of the surface intercepting the sound. The intensity I of a sound wave is given by

$$I = \frac{1}{2} v \rho \omega^2 s_m^2.$$









Variation of Intensity with Distance

Consider a sound source S that emits sound isotropically (with equal intensity in all directions), as shown in the figure.

Imagine a sphere of radius r, centered on the source. All  $\downarrow$  the energy emitted by the source must pass through the surface of the sphere.

Therefore, the time rate at which energy is transferred through the surface must equal the time rate at which energy is emitted by the source (the power  $P_s$  of the source).



Variation of Intensity with Distance

The intensity I at a sphere of radius r must then by

$$I = \frac{P_s}{A_{\rm sphere}} = \frac{P_s}{4\pi r^2}.$$



#### CHECKPOINT 2

The figure indicates three small patches 1, 2, and 3 that lie on the surfaces of two imaginary spheres; the spheres are centered on an isotropic point source S of sound. The rates at which energy is transmitted through the three patches by the sound waves are equal. Rank the patches according to (a) the intensity of the sound on them and (b) their area, greatest first.

(a) 1 & 2 tie, 3. (b) 3, 1& 2 tie.



$$I = \frac{P_s}{4\pi r^2}$$
$$P = IA$$

**Example 3**: A firecracker, emitting a pulse of sound that travels radially outward from the explosion center. The power of this acoustic emission is  $P_s = 3.5 \times 10^4$  W.

(a) What is the intensity I of the sound when it reaches a distance r = 100 m from the firecracker?

$$I = \frac{P_s}{4\pi r^2} = \frac{3.5 \times 10^4 \text{ W}}{4\pi (100 \text{ m})^2} = 0.279 \text{ W/m}^2.$$



(b) At what time rate  $P_{he}$  is sound energy intercepted by a human ear, aimed at the firecracker and located a distance r = 100 m from the firecracker? The radius of the human ear canal is 0.35 mm.

 $A_{he} = \pi r^2 = \pi (0.35 \times 10^{-3} \text{ m})^2 = 3.85 \times 10^{-7} \text{ m}^2.$  $P_{he} = IA_{he} = (0.279 \text{ W/m}^2)(3.85 \times 10^{-7} \text{ m}^2)$  $= 1.1 \times 10^{-7} \text{ W}.$ 



**Example 4**: An electric spark jumps along a straight line of length L = 10 m, emitting a pulse of sound that travels radially outward from the spark. (The spark is said to be a line source of sound.) The power of this acoustic emission is  $P_s = 1.6 \times 10^4$  W.

(a) What is the intensity I of the sound when it reaches a distance r = 12 m from the spark?

The radial sound pulse passes through a cylindrical surface of area  $A = 2\pi rL$ . The intensity is then

$$I = \frac{P_s}{2\pi rL} = \frac{1.6 \times 10^4 \text{ W}}{2\pi (12 \text{ m})(10 \text{ m})} = 21 \text{ W/m}^2.$$



(b) At what time rate  $P_d$  is sound energy intercepted by an acoustic detector of area  $A_d = 2.0 \text{ cm}^2$ , aimed at the spark and located a distance r = 12 m from the spark?

$$P_d = IA_d = (21 \text{ W/m}^2)(2.0 \times 10^{-4} \text{ m}^2) = 4.2 \text{ mW}.$$

(b) At what time rate  $P_d$  is sound energy intercepted by an acoustic detector of area  $A_d = 2.0 \text{ cm}^2$ , aimed at the spark and located a distance r = 12 m from the spark?

 $P_d = IA_d = (21 \text{ W/m}^2)(2.0 \times 10^{-4} \text{ m}^2) = 4.2 \text{ mW}.$ 

The **anechoic chamber** at Orfield Laboratories. It holds the Guinness World Record for the world's quietest place.

 $\beta = -9.4 \text{ dB}$ 



#### The Decibel Scale

It is much more convenient to speak of **sound level**  $\beta$  instead of sound intensity *I*. The sound level is defined as

$$\beta = (10 \text{ dB}) \log \frac{I}{I_0}.$$

Here dB is the abbreviation of **decibel** and  $I_0$  is a standard reference intensity (=  $10^{-12}$  W /m<sup>2</sup>), chosen near the lowest limit of human range of hearing.

 $\beta = 60 \text{ dB}$  corresponds to an intensity that is  $10^6$  times the standard reference level.

#### Some Sound Levels (dB)

Hearing threshold	0
Rustle of leaves	10
Conversation	60
Rock concert	110
Pain threshold	120
Jet engine	130

**Example 4**: If earplugs decreases the sound level of some source by 20 dB, what is the ratio of the final intensity  $I_f$  of the sound waves to their initial intensity  $I_i$ ?

$$\beta_f - \beta_i = (10 \text{ dB}) \left( \log \frac{I_f}{I_0} - \log \frac{I_i}{I_0} \right) = (10 \text{ dB}) \log \frac{I_f}{I_0} \frac{I_0}{I_i} = (10 \text{ dB}) \log \frac{I_f}{I_i}.$$

Solving for  $I_f/I_i$  we obtain

$$\frac{I_f}{I_i} = 10^{\frac{\beta_f - \beta_i}{10 \text{ dB}}}.$$

Using  $\beta_f - \beta_i = -20 \text{ dB}$  yields

$$\frac{I_f}{I_i} = 10^{\frac{-20 \text{ dB}}{10 \text{ dB}}} = 10^{-2}.$$

Musical sounds can be set by oscillating strings, membranes, air columns and so on.

We have seen in Ch. 16 that setting up standing waves on a string make it oscillate with a large sustained amplitude.

We can set up standing sound waves in an air-filled pipe, with one or two open ends. The waves get reflected at both ends of the pipe, whether closed or open! If the wavelength is suitably matched to the length of the pipe, the superposition of the waves in the pipe sets up a standing wave pattern.

The advantage of such a standing wave is that the air in the pipe oscillates with large, sustained amplitude, emitting sound at any open end at the same frequency of the standing wave in the pipe.

Standing waves in a pipe are similar to those in a string. There is an antinode at an open end of a pipe and a node at a closed end.

The simplest standing wave pattern (the fundamental mode or first harmonic), that can be set in a pipe with **two open ends**, is shown in the figure.

The wavelength  $\lambda$  is twice the pipe's length L, or  $\lambda = 2L$ .



The second harmonic requires that  $\lambda = L$ . The third harmonic requires that  $\lambda = 2L/3$ .

Generally, the resonant frequency for a pipe of length L with two open ends corresponds to  $\frac{1}{n} = n = 1$ 

$$\lambda = \frac{2L}{n}, \qquad n = 1, 2, 3, \dots$$

The corresponding resonant frequencies are

$$f = \frac{v}{\lambda} = n \frac{v}{2L}, \qquad n = 1, 2, 3, \dots$$



For a pipe with **one open end**, there must be an antinode at the open end and a node at the closed end.

The simplest standing wave pattern requires that  $L = \lambda/4$ , or  $\lambda = 4L$ .

The next simplest pattern requires that  $L = 3\lambda/4$  or  $\lambda = 4L/3$ , and so on.



Generally, the resonant frequency for a pipe of length L with one open end corresponds to wavelengths

$$\lambda = \frac{4L}{n}$$
,  $n = 1, 3, 5, ... (odd n)$ 

The corresponding resonant frequencies are

$$f = \frac{v}{\lambda} = n \frac{v}{4L}, \qquad n = 1, 3, 5, \dots \text{ (odd n)}$$



# CHECKPOINT 3

Pipe A, with length L, and pipe B, with length 2L, both have two open ends. Which harmonic of pipe B has the same frequency as the fundamental of pipe A?

#### The second harmonic

$$f = n \frac{v}{2L}$$

$$f_{A} = n_{A} \frac{v}{2L}; \qquad f_{A1} = \frac{v}{2L}$$

$$f_{B} = n \frac{v}{2(2L)} = n_{B} \frac{v}{4L}; \qquad f_{B2} = 2 \frac{v}{4L} = \frac{v}{2L}$$

**Example 5**: Pipe A is open at both ends and has length  $L_A = 0.343$  m. We want to place it near three other pipes in which standing waves have been set up, so that the sound can set up a standing wave in pipe A. Those other three pipes are each closed at one end and have lengths  $L_B = 0.500L_A$ ,  $L_C = 0.250L_A$ , and  $L_D = 2.00L_A$ . For each of these three pipes, which of their harmonics can excite a harmonic in pipe A? What frequency do you hear from the tube?

Pipe A:

$$f_A = \frac{n_A v}{2L_A} = n_A \frac{343 \text{ m/s}}{2(0.3430 \text{ m})} = n_A (500 \text{ Hz}), \qquad n_A = 1,2,3,...$$

Pipe B:

$$f_B = \frac{n_B v}{4L_B} = n_B \frac{343 \text{ m/s}}{4(0.500L_A)} = n_B \frac{343 \text{ m/s}}{4(0.500)(0.3430 \text{ m})} = n_B (500 \text{ Hz}),$$
$$n_B = 1,3,5, \dots$$

Thus, 
$$f_A = f_B$$
 for  $n_A = n_B$ , with  $n_B = 1,3,5,...$   
Pipe C:

$$f_C = \frac{n_C v}{4L_C} = n_C \frac{343 \text{ m/s}}{4(0.250L_A)} = n_C \frac{343 \text{ m/s}}{4(0.250)(0.3430 \text{ m})} = n_C (1000 \text{ Hz}),$$
$$n_C = 1,3,5, \dots$$

Thus, 
$$f_A = f_C$$
 for  $n_A = 2n_C$ , with  $n_C = 1,3,5,...$ 

Pipe D:

$$f_{D} = \frac{n_{D}v}{4L_{D}} = n_{D}\frac{343 \text{ m/s}}{4(2L_{A})}$$
  
=  $n_{D}\frac{343 \text{ m/s}}{4(2)(0.3430 \text{ m})}$   
=  $n_{D}(125 \text{ Hz})$   $n_{D} = 1,3,5,...$   
Thus,  $f_{A} = f_{D}$  for  $n_{A} = \frac{1}{2}n_{D}$ , which is not possible since  $n_{D}$  is odd.



**Example 6**: Weak background noises from a room set up the fundamental standing wave in a cardboard tube of length L = 67.0 cm with two open ends. Assume that the speed of sound in the air within the tube is 343 m/s.

(a) What frequency do you hear from the tube?

$$f_1 = \frac{v}{2L} = \frac{343 \text{ m/s}}{2(0.670 \text{ m})} = 256 \text{ Hz}.$$

(b) If you jam your ear against one end of the tube, what fundamental frequency do you hear from the tube?

The tube now has one open end. Thus

$$f_1 = \frac{v}{4L} = \frac{343 \text{ m/s}}{4(0.670 \text{ m})} = 128 \text{ Hz.}$$

https://www.youtube.com/watch?v=a3RfULw7aAY

Consider a source S of sound and a detector D that are in motion relative to air, with speeds  $v_S$  and  $v_D$ , respectively. The emitted frequency f and detected frequency f' are related by

$$f' = f \frac{v \pm v_D}{v \pm v_S},$$

where v is the speed of sound through air.

Sign Rule:

- When the source or detector moves toward the other, the sign on its speed must give an upward shift in frequency.
- When the source or detector moves away from the other, the sign on its speed must give a downward shift in frequency.

#### CHECKPOINT 4

The figure indicates the directions of motion of a sound source and a detector for six situations in stationary air. For each situation, is the detected frequency greater than or less than the emitted frequency, or can't we tell without more information about the actual speeds?

	Source	Detector		Source	Detector
(a)	$\longrightarrow$	• 0 speed	(d)	←	←
(b)	←───	• 0 speed	(e)	$\longrightarrow$	←───
(c)	$\longrightarrow$	$\longrightarrow$	( <i>f</i> )	←	$\longrightarrow$

(a) 
$$f' = f \frac{v}{v - v_S}$$
, greater  
(b)  $f' = f \frac{v}{v + v_S}$ , less  
(c)  $f' = f \frac{v - v_D}{v - v_S}$ , can't tell  
(d)  $f' = f \frac{v + v_D}{v + v_S}$ , can't tell  
(e)  $f' = f \frac{v + v_D}{v - v_S}$ , greater  
(f)  $f' = f \frac{v - v_D}{v + v_S}$ , less

**Example 5**: Bats navigate and search out prey by emitting, and then detecting reflections of, ultrasonic waves, which are sound waves with frequencies greater than can be heard by a human. Suppose a bat emits ultrasound at frequency  $f_{be} = 82.52$  kHz while flying with velocity  $\vec{v}_b = (9.00 \text{ m/s})\hat{i}$  as it chases a moth that flies with velocity  $\vec{v}_m = (8.00 \text{ m/s})\hat{i}$ . What frequency  $f_{md}$  does the moth detect?

The source (bat) is moving *toward* the detector (moth). The detector (moth) is moving *away* from the source (bat). Thus,

$$f_{md} = f_{be} \frac{v - v_D}{v - v_S} = (82.52 \text{ kHz}) \frac{343 \text{ m/s} - 8.00 \text{ m/s}}{343 \text{ m/s} - 9.00 \text{ m/s}} = 82.8 \text{ kHz}.$$

What frequency  $f_{bd}$  does the bat detect in the returning echo from the moth? The source (moth) is moving *away* from the detector (bat). The detector (bat) is moving *toward* the source (moth). Thus,

$$f_{bd} = f_{md} \frac{v + v_D}{v + v_S} = (82.77 \text{ kHz}) \frac{343 \text{ m/s} + 9.00 \text{ m/s}}{343 \text{ m/s} + 8.00 \text{ m/s}} = 83.0 \text{ kHz}$$