

In what follows we consider Laplace integrals of the type in which $\phi(t) = t^\lambda g(t)$, where $g(t)$ possesses Taylor series expansion about $t=0$ with $g(0) \neq 0$ and λ is real and $\lambda > -1$. This ensure convergence of the integral at $t=0$. If $g(t) = 0$ having a zero of order r at $t=0$ then it has a factor t^r and we re-define $g(t)$ such that now $g(t) \neq 0$ at $t=0$ while $g t^r$ is absorbed in t^λ .

Watson Lemma : $\int_0^\infty e^{-xt} t^\lambda g(t) dt \sim \sum_{n=0}^{\infty} \frac{g^{(n)}(0) \Gamma(\lambda+n+1)}{n! x^{\lambda+n+1}}$ as $x \rightarrow \infty$.

(x real, +ve, T some +ve number)
Proof : $|g(t)| < k e^{ct}$ as $x \rightarrow \infty$.

Since $g(t)$ has a Taylor series expansion about $t=0$,

$$g(t) = \sum_{n=0}^{\infty} \frac{g^{(n)}(0)}{n!} t^n = \sum_{n=0}^{\infty} a_n t^n = \sum_{n=0}^N a_n t^n + r_N(t) \quad (1)$$

where $|r_N(t)| < L t^{N+1}$ for $|t| < R$

and $a_n = \frac{g^{(n)}(0)}{n!}$.

We shall consider the radius of convergence $R > T$ but this is not necessary. In the integral, we use (1), to get

$$f(x) = \int_0^T e^{-xt} t^\lambda \sum_{n=0}^N a_n t^n + \int_0^T e^{-xt} t^\lambda r_N(t) dt. \quad (2)$$

We now concentrate on the second integral on the R.H.S

$$\left| \int_0^T e^{-xt} t^\lambda r_N(t) dt \right| < L \int_0^T e^{-xt} t^{\lambda+N+1} dt \quad (3)$$

(Wat2)

Put $\tau = xt$ ($d\tau = x dt$) on R.H.S

$$\begin{aligned}
 \text{R.H.S.} &= \int_0^T e^{-xt} t^{\lambda+N+1} dt = \frac{-(\lambda+N+2)}{x} \int_0^{xT} e^{-\tau} \tau^{\lambda+N+1} d\tau \\
 &= \frac{-(\lambda+N+2)}{x} \left[\int_0^{\infty} e^{-\tau} \tau^{\lambda+N+1} d\tau - \int_{xT}^{\infty} e^{-\tau} \tau^{\lambda+N+1} d\tau \right] \\
 &= \frac{-(\lambda+N+2)}{x} \Gamma(\lambda+N+2) - \underbrace{\frac{-(\lambda+N+2)}{x} \int_{xT}^{\infty} e^{-\tau} \tau^{\lambda+N+1} d\tau}_{\text{second integral}}
 \end{aligned}$$

Again consider second integral on R.H.S. Put

 $\tau = xT(1+u)$ ($d\tau = xT du$)

$$\begin{aligned}
 \frac{-(\lambda+N+2)}{x} \int_{xT}^{\infty} e^{-\tau} \tau^{\lambda+N+1} d\tau &= T^{\lambda+N+2} e^{-xT} \int_0^{\infty} e^{-xTu} (1+u)^{\lambda+N+1} du \\
 &< T^{\lambda+N+2} e^{-xT} \int_0^{\infty} e^{-xTu} e^{(\lambda+N+1)u} du \\
 &\quad \left(\text{using } (1+u)^a < e^{au}, a > 0, u > 0 \right) \\
 &= T^{\lambda+N+2} e^{-xT} \frac{1}{xT - (\lambda+N+1)} \\
 &= T^{\lambda+N+2} e^{-xT} \frac{1}{xT} \left(1 + \frac{\lambda+N+1}{xT} + \dots \right) \\
 &\sim T^{\lambda+N+1} \frac{e^{-xT}}{x},
 \end{aligned}$$

Going back to $\int_0^T$ R.H.S we see

$$\begin{aligned}
 \int_0^T e^{-xt} t^{\lambda+N+1} dt &= \frac{-(\lambda+N+2)}{x} \Gamma(\lambda+N+2) + o(e^{-xT}) \\
 &= \frac{-(\lambda+N+2)}{x} \Gamma(\lambda+N+2) + o(x^{-(\lambda+N+2)}) \\
 &\quad \text{as } x \rightarrow \infty.
 \end{aligned}$$

Using this estimate in (3)

$$\int_0^T e^{-xt} t^\lambda r_N(t) dt = O(x^{-(\lambda+N+2)}) \text{ as } x \rightarrow \infty.$$

Thus (2) becomes

$$\begin{aligned} \text{Hence } \int_0^T e^{-xt} t^\lambda \sum_{n=0}^{\infty} a_n t^n dt &= \sum_{n=0}^N a_n \left[\int_0^{\infty} - \int_1^{\infty} \right] e^{-xt} t^{\lambda+n} dt \\ &= \sum_{n=0}^N a_n \Gamma(\lambda+n+1) x^{-(\lambda+n+1)} + O(x^{-(\lambda+N+1)}) \\ &\quad \text{as } x \rightarrow \infty. \end{aligned}$$

$$\text{Thus } f(x) = \sum_{n=0}^N a_n \Gamma(\lambda+n+1) x^{-(\lambda+n+1)} + O(x^{-(\lambda+N+2)})$$

which gives the required asymptotic formula

$$\int_0^T e^{-xt} t^\lambda g(t) dt \sim \sum_{n=0}^{\infty} \frac{g^{(n)}(0) \Gamma(\lambda+n+1)}{n! x^{\lambda+n+1}} \text{ as } x \rightarrow \infty.$$

Remark: This result can also be shown for complex z as long as we take $|\arg z| < \pi/2$. Proof can be seen in Copson (1965: Asymptotic expansions, Cambridge University Press)

(2) The upper limit T can be replaced by ∞ without requiring the radius of convergence R to be infinite.