#### TRIDIAGONAL PHYSICS

Part I: 20 min. Introduction Formulation Applications & Examples Perspectives Part II: 15 min., Dr. Abdelmonem The J-matrix method Resonance search

#### Introduction

Tridiagonal Physics is the Physics obtained from the analysis of the matrix representation of relevant operators (e.g., the Hamiltonian) in a basis in which it is tridiagonal.

Systematic program started in 2000

- Multi-disciplinary program
- Dhahran group + national/international membership, collaboration, and support
- Rich agenda: long list of significant problems

#### Formulation

Consider H : diagonal rep.  $H\phi_n = E_n\phi_n$ too restrictive: discrete spectrum, small set of potentials Relax constraint:  $H\phi_n \sim \phi_n + \phi_{n-1} + \phi_{n+1}$  $\begin{bmatrix} a_0 & b_0 \end{bmatrix}$  $\left|\psi\right\rangle = \sum f_{n} \left|\phi_{n}\right\rangle$  $b_0 \quad a_1 \quad b_1 \qquad \mathbf{0}$  $H |\psi\rangle = x |\psi\rangle \qquad (1) \longrightarrow$  $H = \begin{bmatrix} b_{1} & a_{2} & b_{2} \\ b_{1} & a_{2} & b_{2} \\ b_{2} & \times & \times \end{bmatrix} \qquad H = \begin{bmatrix} \psi & f_{1} & \psi & f_{1} \\ f_{1} & \phi & f_{2} & 0 \\ f_{2} & f_{1} & f_{2} & 0 \\ f_{1} & f_{2} & f_{2} & 0 \\ f_{2} & f_{1} & f_{2} & f_{2} & 0 \\ f_{1} & f_{2} & f_{2} & f_{2} & 0 \\ f_{2} & f_{2} & f_{2} & f_{2} & f_{2} & 0 \\ f_{2} & f_{2} & f_{2} & f_{2} & f_{2} & f_{2} & f_{2} \\ f_{2} & f_{2} \\ f_{2} & f_$  $xf_0 = a_0f_0 + b_0f_1$   $\langle \phi_n | \phi_m \rangle = \delta_{nm}$  $0 \times \times \times$  $|\psi\rangle \Leftrightarrow \{f_n\}$  (1)  $\Leftrightarrow$  (2)

Spectrum  $\{x\}$  includes discrete as well as continuous and larger class of potentials

 $xf_{n}(x) = a_{n}f_{n}(x) + b_{n-1}f_{n-1}(x) + b_{n}f_{n+1}(x) \qquad \int_{x}^{x+} \rho(x)f_{n}(x)f_{m}(x)dx = \delta_{nm}$ Normalization:  $f_{0} = 1 \longrightarrow f_{n}$  polynomial in x of degree n

Examples of such orthogonal polynomials:

- Hermite:  $\frac{1}{\sqrt{2^n n!}} H_n(x), \quad x \in [-\infty, +\infty], \quad a_n = 0, b_n = \sqrt{\frac{n+1}{2}}$
- Cheybechev:  $T_n(x)$ ,  $x \in [-1, +1]$ ,  $a_n = 0$ ,  $b_n = \frac{1}{2}$

• Laguerre:  $\sqrt{\frac{\Gamma(n+1)\Gamma(\nu+1)}{\Gamma(n+\nu+1)}}L_n^{\nu}(x)$ ,  $x \in [0, +\infty]$ ,  $a_n = 2n + \nu + 1$ ,  $b_n = -\sqrt{(n+1)(n+\nu+1)}$ 

 Associated tools: Lanczos algorithm, Gauss quadratures, continued fractions, Pade approximations, ...etc.

• In the literature: Toda lattices, tight-binding models, recursion methods, 1D chain models, ...etc.

### Applications & Examples The J-matrix method

- Algebraic method of quantum scattering (1974, 1975)
- Atomic, molecular, nuclear, chemical
- Accuracy and convergence
- Multi-channel (rigorous formulation, 1997)
- Relativistic extension (2000)
- Book, Springer (2007)
- Active Groups: Russia, Saudi Arabia, North Europe, Australia, China, Brazil.
- Dr. M. S. Abdelmonem (Part II)

#### Discrete & Continuous Spectrum (all energies)

 $H\phi_n = \overline{E}_n\phi_n$  $E \phi_n = a_n \phi_n + b_{n-1} \phi_{n-1} + b_n \phi_{n+1}$ • Larger class of solvable potentials: Ann. Phys. 317, 152 (2005) • Example: analytic solution for a new noncentral potential  $V(r,\theta) = \alpha \frac{\cos \theta}{r^2}$ 

J. Phys. A 38, 3409 (2005)

#### Spectral Density

Resolvent operator:  $G_{nm}(z) = (H - z)_{nm}^{-1}$ Project density:  $\rho_n(x) = \frac{1}{2\pi i} [G_{nn}(x+i\ 0) - G_{nn}(x-i\ 0)]$ 

*H* finite  $\rightarrow G_{nn}$  non-analytic: interleaved poles & zeros Three approximation methods:  $(a_0, b_0)$ 

- Analytic continuation
- Dispersion correction
- Stieltjes imaging

Phys. Rev. A 62, 052103 (2005)

 $\lim_{n \to \infty} \{a_n, b_n\} = \{a_m, b_m\}_{m=1}^K \longrightarrow K \text{-band density with } K-1 \text{ energy gaps}$ 

• Impurity at site k:  $a_k \rightarrow a_k + \mu$  and/or  $b_k \rightarrow \gamma b_k$ Density deformation: new density in terms of original one and deformation parameters  $\mu$  and  $\gamma$ .

Phys. Lett. A (2007), in production

• Density associated with nonlinear bifurcating map that exhibit period doubling cascade:

No. bands = half of the number of fixed points of the map (orbits period or bifurcation branches)

J. Phys. A 39, 6851 (2006)

Deformation of orthogonal polynomials

 $a_k \rightarrow a'_k \quad b_k \rightarrow b'_k \neq 0$  for a given integer k  $\rho(x) \rightarrow \rho'(x) \qquad P_n(x) \rightarrow P'_n(x)$  $\int_{x_-}^{x_+} \rho'(x) P'_n(x) P'_m(x) dx = \delta_{nm}$ 

J. Phys. A 35, 9071 (2002)

## Study of resonance in the complex charge plane

Complex *E*-plane: for a given angular momentum and charge: bound states & resonances, complex scaling

Complex  $\ell$ -plane: for a given *E* and *Z*: Regge poles (  $\operatorname{Re} \ell = 0, 1, ...$ ), Regge trajectories (vary *E*)

**Complex Z-plane**:  $H - E = H_0 + V + \frac{Z}{r} - E$   $r(E - H_0 - V)|\psi\rangle = Z|\psi\rangle$  Tridiagonal rep.  $\rightarrow$  straight-forward application of complex scaling Scattering for a given  $E \& \ell$  but for all Z Regge-like poles (Re  $Z = 0, \pm 1, \pm 2, ...$ ) Regge-like trajectories (vary E) J. Phys. A 37, 5863 (2004) Evaluation of new integrals involving orthogonal polynomials

 $\int_{0}^{\infty} x^{\nu} e^{-x/2} J_{\nu}(\mu x) L_{n}^{2\nu}(x) dx =$   $2^{\nu} \Gamma\left(\nu + \frac{1}{2}\right) \frac{1}{\sqrt{\pi\mu}} (\sin\theta)^{\nu + \frac{1}{2}} C_{n}^{\nu + \frac{1}{2}} (\cos\theta)$ 

Where  $\mu$  and  $\nu$  are real parameters such  $\mu \ge 0$ ,  $\nu > -\frac{1}{2}$ .  $\cos \theta = \frac{\mu^2 - \frac{1}{4}}{\mu^2 + \frac{1}{4}}$  and  $C_n^{\lambda}(x)$  is the Gegenbauer polynomial.

Appl. Math. Lett. 20, 38 (2007)

#### Perspectives

#### Group's work plan:

Accurate evaluation of bound states & resonances
 Solutions of the Dirac equation with non-central potential for all energies

- J-matrix inspired quadrature & matrix tridiagonalization method
- Spectral decomposition of the Coulomb wave function
  ...etc.

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Thank you