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Note

A note on blow up of solutions of a quasilinear heat equation with vanishing initial energy

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Abstract

In this work we consider an initial boundary value problem related to the equation

$$u_t - \operatorname{div}(|\nabla u|^{m-2}\nabla u) = f(u)$$

and prove, under suitable conditions on f, a blow up result for solutions with vanishing or negative initial energy.

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1. Introduction

In this paper we are concerned with the finite time blow up of solutions for the initial boundary value problem

$$u_t - \operatorname{div}(|\nabla u|^{m-2}\nabla u) = f(u), \quad x \in \Omega, \ t > 0,$$

$$u(x,t) = 0, \quad x \in \partial\Omega, \ t \ge 0,$$

$$u(x,0) = u_0(x), \quad x \in \Omega,$$
(1.1)

where m > 2, f is a continuous function, and Ω is a bounded domain of \mathbb{R}^n $(n \ge 1)$, with a smooth boundary $\partial \Omega$.

In 1993, Junning [2] studied (1.1) and established a global existence result for f depending on u as well as on ∇u . He also proved a nonglobal existence result for (1.1) under the condition

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$$\frac{1}{m} \int_{\Omega} \left| \nabla u_0(x) \right|^m dx - \int_{\Omega} F\left(u_0(x)\right) dx \leqslant -\frac{4(m-1)}{mT(m-2)^2} \int_{\Omega} u_0^2(x) dx,$$
(1.2)

where $F(u) = \int_0^u f(s) ds$. More precisely, he showed that if there exists T > 0, for which (1.2) holds, then the solution blows up in a time less than T. This type of results have been extensively generalized and improved by Levine et al. in [3], where the authors proved some global, as well as nonglobal, existence theorems. Their result, when applied to problem (1.1), requires that

$$\frac{1}{m} \int_{\Omega} \left| \nabla u_0(x) \right|^m dx - \int_{\Omega} F\left(u_0(x)\right) dx < 0. \tag{1.3}$$

We note that the inequality (1.3) implies (1.2). In 1999, Erdem [1] discussed the initial Dirichlet-type boundary problem for

$$u_{t} - \sum_{i=1}^{n} \frac{\partial}{\partial x_{i}} \left(\left(d + |\nabla u|^{m-2} \right) \frac{\partial u}{\partial x_{i}} \right) + g(u, \nabla u) = f(u),$$

$$x \in \Omega, \ t > 0. \tag{1.4}$$

He established a blow up result, under a condition similar to (1.3) and another one on the growth of g.

Concerning global existence, Nakao and Chen [5] studied the following problem:

$$u_t - \operatorname{div}(\sigma(|\nabla u|^2)\nabla u) + b(u)\nabla u = 0, \quad x \in \Omega, \ t > 0,$$

$$u(x,t) = 0, \quad x \in \partial\Omega, \ t \geqslant 0,$$

$$u(x,0) = u_0(x), \quad x \in \Omega,$$
(1.5)

where $\sigma(v)$ behaves like $|v|^m$, $m \ge 0$, and $|b(u)| \le k_0 |u|^\beta$, $k_0 > 0$, $\beta \ge 0$. He proved global existence, derived precise estimates for $\nabla u(t)$, and showed that solutions decay as $t \to \infty$. His work improves an earlier one by Nakao and Ohara [4], in which he considered (1.5) with $b \equiv 0$.

It is also worth mentioning that Nakao and Ohara [6] considered the periodic solutions of (1.5), with the last term replaced by g(x,u)-f(x,t). He showed that these periodic solutions belong to $L^{\infty}(\omega,W^{1,\infty}(\Omega))$ and gave a bound of $\|\nabla u(t)\|_{\infty}$ under certain geometric conditions on $\partial\Omega$.

Here we show that the blow up can be obtained even for vanishing energy. More precisely, we will get a blow up under the condition

$$\frac{1}{m} \int_{\Omega} \left| \nabla u_0(x) \right|^m dx - \int_{\Omega} F\left(u_0(x)\right) dx \leqslant 0. \tag{1.6}$$

To make this paper self-contained we state, without proof, the local existence result of [2].

Proposition. Let f be in $C(\mathbb{R})$ satisfying

$$\left| f(u) \right| \leqslant g(u) \tag{1.7}$$

for g a C^1 function. Then for any $u_0 \in L^{\infty}(\Omega) \cap H_0^m(\Omega)$, the problem (1.1) has a solution

$$u \in L^{\infty}(\Omega \times (0, T)) \cap L^{m}((0, T); H_{0}^{m}(\Omega)),$$

$$u_{t} \in L^{2}(\Omega \times (0, T)).$$
(1.8)

2. Blow up

In this section we state and prove our main result.

Theorem. Let f be in $C(\mathbb{R})$ satisfying (1.7) and

$$pF(u) \leqslant uf(u), \quad p > m > 2. \tag{2.1}$$

Then for any nonzero $u_0 \in L^{\infty}(\Omega) \cap H_0^m(\Omega)$ satisfying (1.6), the solution (1.8) blows up in finite time.

Remark. An example of a function f satisfying (2.1) is $f(s) = |s|^{p-2}s$, for p > m > 2. This shows that, in a sense, the source has to dominate the m-Laplacian term.

Proof. We define

$$H(t) = \int_{\Omega} F(u(x,t)) dx - \frac{1}{m} \int_{\Omega} |\nabla u(x,t)|^m dx.$$

By using (1.1), we easily arrive at

$$H'(t) = \int_{\Omega} u_t^2(x, t) \, dx \geqslant 0;$$

hence $H(t) \ge H(0) \ge 0$, by virtue of (1.6). We then set

$$L(t) = \frac{1}{2} \int_{\Omega} u^2(x, t) dx$$

and differentiate L to get

$$L'(t) = \int_{\Omega} u u_t(x, t) \, dx \geqslant \int_{\Omega} u \left[\operatorname{div} \left(|\nabla u|^{m-2} \nabla u \right) + f(u) \right] (x, t) \, dx$$

$$\geq pH(t) + \left(\frac{p}{m} - 1\right) \int_{\Omega} |\nabla u|^{m}(x, t) dx$$

$$\geq \left(\frac{p}{m} - 1\right) \left[H(t) + \|\nabla u\|_{m}^{m}\right] \geq 0. \tag{2.2}$$

Next we estimate $L^{m/2}(t)$:

$$L^{m/2}(t) \leqslant C \|u\|_m^m \leqslant C \|\nabla u\|_m^m$$

by Poincare's inequality and the embedding of the L^q spaces. Here C is a constant depending on Ω and m only. Therefore we have

$$L^{m/2}(t) \leqslant C[H(t) + \|\nabla u\|_m^m].$$
 (2.3)

By combining (2.2) and (2.3) we have

$$L'(t) \geqslant \gamma L^{m/2}(t), \tag{2.4}$$

where $\gamma = (p - m)/Cm$. A direct integration of (2.4) then yields

$$L^{m/2-1}(t) \geqslant \frac{1}{L^{1-m/2}(0) - \gamma t}.$$

Therefore *L* blows up in a time $t^* \leq 1/\gamma L^{(m/2)-1}(0)$.

Corollary. *If there exists* $t_0 \ge 0$, *for which*

$$\frac{1}{m} \int_{\Omega} \left| \nabla u(x, t_0) \right|^m dx - \int_{\Omega} F\left(u(x, t_0)\right) dx = 0,$$

then the solution (1.8) either remains equal to zero for all time $t \ge t_0$ or blows up in finite time $t^* > t_0$.

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