

13. *Torque in a Couette viscometer*—M. Fig. P6.13 shows the horizontal cross section of a concentric cylinder or “Couette” viscometer, which is an apparatus for determining the viscosity μ of the fluid that is placed between the two vertical cylinders. The inner and outer cylinders have radii of r_1 and r_2 , respectively. If the inner cylinder is rotated with a steady angular velocity ω , and the outer cylinder is stationary, derive an expression for v_θ (the θ velocity component) as a function of radial location r .

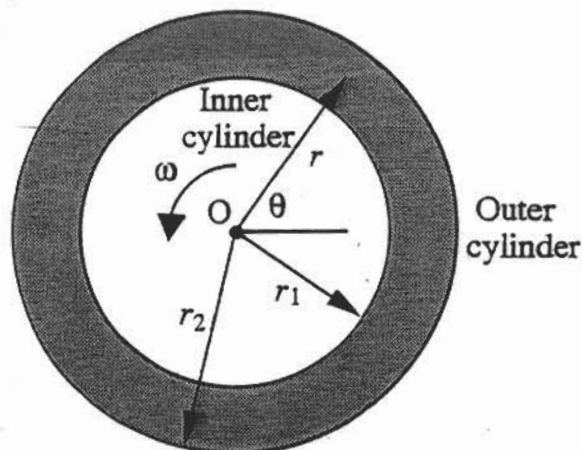


Fig. P6.13 Section of a Couette viscometer.

If, further, the torque required to rotate the *inner* cylinder is found to be T per unit length of the cylinder, derive an expression whereby the unknown viscosity μ can be determined, in terms of T , ω , r_1 , and r_2 . *Hint*: you will need to consider one of the shear stresses given in Table 5.8.

$$T = \text{Force} \times r_1 \quad (1)$$

$$g_z = -g, \quad g_\theta = g_r = 0$$

To oppose the effect of the shear stress

$$\text{Force} = -\tau_{r\theta} \Big|_{r=r_1} \times 2\pi r_1 L \quad (2)$$

the surface on which the sheare stress acts

direction of the shear stress

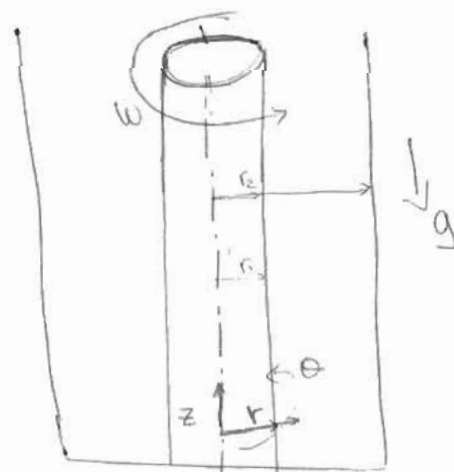


Table 5.8

$$\tau_{r\theta} = \mu \left[r \frac{\partial}{\partial r} \left(\frac{u_\theta}{r} \right) + \frac{1}{r} \frac{\partial u_r}{\partial \theta} \right] \quad (3)$$

Assumptions

- 1) Steady State
- 2) Newtonian fluid (μ is constant)
- 3) Incompressible fluid (ρ is constant)
- 4) $u_z = u_r = 0$ Velocity in the θ direction
- 5) Axisymmetric problem (all the variables are not $f(\theta)$)

Continuity Equation

ρ is constant, $u_r = u_z = 0$

$$\therefore \frac{\partial u_\theta}{\partial \theta} = 0$$

$$\therefore u_\theta = f(r)$$

Momentum balance

$$r\text{-direction: } \frac{-\rho u_\theta^2}{r} = -\frac{\partial P}{\partial r} \quad \left(\frac{\partial P}{\partial \theta} = 0 \right) \quad (4)$$

$$\theta\text{-direction: } \mu \left[\frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial (u_\theta r)}{\partial r} \right) \right] = 0 \quad (5)$$

$$z\text{-direction: } 0 = -\frac{\partial P}{\partial z} + \rho g_z \Rightarrow g_z = -g \quad (6)$$

$$P = f(r, z)$$

equation (6)

$$\frac{\partial P}{\partial z} = -\rho g$$

Integration gives $P = -\rho g z + f_1(r)$

Equation (4)

$$u_{\theta} = f(r)$$

$$P = \int \frac{u_{\theta}^2}{r} dr + f_2(z)$$

$$\therefore f_1(r) = \int \frac{u_{\theta}^2}{r} dr, \quad f_2(z) = -\rho g z$$

equation (5)

$$\frac{d}{dr} \left(\frac{1}{r} \frac{d(r u_{\theta})}{dr} \right) = 0$$

$$\text{Integration} \Rightarrow \frac{1}{r} \frac{d(r u_{\theta})}{dr} = C_1$$

$$\text{or } \frac{d(r u_{\theta})}{dr} = C_1 r$$

$$\text{Integration} \Rightarrow r u_{\theta} = C_1 \frac{r^2}{2} + C_2$$

$$\text{or } u_{\theta} = C_1 \frac{r}{2} + \frac{C_2}{r}$$

Boundary Condition

$$r = r_1$$

$$r = r_2$$

$$u_{\theta} = \omega r_1$$

$$u_{\theta} = 0$$

angular velocity $\frac{1}{s}$

No slip condition

$$\therefore \omega r_1 = C_1 \frac{r_1}{2} + \frac{C_2}{r_1}$$

$$0 = C_1 \frac{r_2}{2} + \frac{C_2}{r_2}$$

Solving for C_1 & C_2 and substituting in the expression of u_{θ} gives

$$u_{\theta} = \frac{r_1 \omega}{r_1 - \frac{r_2^2}{r_1}} r + \frac{r_1 \omega}{\left(\frac{1}{r_1} - \frac{r_1}{r_2^2}\right)} \frac{1}{r}$$

Pressure

$$P = -\rho g z + \int \frac{u_{\theta}^2}{r} dr$$

$$P = -\rho g z + \frac{c_1^2}{4} r^2 + c_1 c_2 \ln r - \frac{c_2^2}{2r^2}$$

$$c_1 = \frac{-2r_1 \omega}{\left(\frac{r_2^2}{r_1} - r_1\right)} \quad , \quad c_2 = \frac{r_1 \omega}{\left(\frac{1}{r_1} - \frac{r_1}{r_2^2}\right)}$$

equation (3)

$$\tau_{r\theta} = \mu \left[r \frac{d}{dr} \left(\frac{u_{\theta}}{r} \right) \right]$$

$$\frac{u_{\theta}}{r} = \frac{c_1}{2} + \frac{c_2}{r^2}$$

$$\frac{d}{dr} \left(\frac{u_{\theta}}{r} \right) = -2c_2/r^3$$

$$\therefore \tau_{r\theta} = -\frac{2\mu c_2}{r^2}$$

equation (2)

$$\text{Force} = \frac{2\mu c_2}{r_1^2} \times 2\pi r_1 L = \frac{4\pi\mu c_2 L}{r_1}$$

Equation (1)

$$T = 4\pi\mu L C_2$$

$$\text{or } T = \frac{4\pi\mu L r_1 \omega}{\left(\frac{1}{r_1} - \frac{r_1}{r_2^2}\right)}$$

Solving for μ gives

$$\mu = \frac{T\left(\frac{1}{r_1} - \frac{r_1}{r_2^2}\right)}{4\pi L r_1 \omega}$$