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King Fahd University of Petroleum & Minerals

Physics Department
Phys630 – "Phase Transition 1" Course
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Exercises # 3

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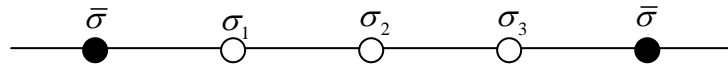
Exercise 8.13

Show that the Bethe- peierls approximation does not predict a phase transition for 1D Ising model with no external field. (As a matter of fact, the Bethe-Peierls approximation gives the exact result for the energy of the 1D Ising model with or with out an external field.)

The solution:

We always look for the solution to Ising model at (external magnetic field $H = 0$) for spins 1 and 3 and a different external field H' for spin 2 .

So, the Hamiltonian is given by



$$E = -H(\sigma_1 + \sigma_3) - H'\sigma_2 - V(\bar{\sigma}\sigma_1 + \sigma_1\sigma_2 + \sigma_2\sigma_3 + \sigma_3\bar{\sigma})$$

$$\Rightarrow Z = \sum_{\sigma_1, \sigma_2, \sigma_3} e^{-\beta E} = \sum_{\sigma_1, \sigma_2, \sigma_3} e^{\beta H(\sigma_1 + \sigma_3) + \beta H'\sigma_2 + \beta V(\bar{\sigma}\sigma_1 + \sigma_1\sigma_2 + \sigma_2\sigma_3 + \sigma_3\bar{\sigma})}$$

$$\Rightarrow Z = \sum_{\sigma_1, \sigma_2, \sigma_3} e^{h(\sigma_1 + \sigma_3) + h'\sigma_2 + v(\bar{\sigma}\sigma_1 + \sigma_1\sigma_2 + \sigma_2\sigma_3 + \sigma_3\bar{\sigma})} \text{ where } \beta H = h, \beta H' = h' \text{ and } \beta V = v$$

$$\Rightarrow Z = \sum_{\sigma_2} e^{h'\sigma_2} \sum_{\sigma_1, \sigma_3} e^{h(\sigma_1 + \sigma_3) + v(\bar{\sigma}\sigma_1 + \sigma_1\sigma_2 + \sigma_2\sigma_3 + \sigma_3\bar{\sigma})} = \sum_{\sigma_2} e^{h'\sigma_2} \sum_{\sigma_1, \sigma_3} e^{\sigma_1(h + v\bar{\sigma} + v\sigma_2) + \sigma_3(h + v\bar{\sigma} + v\sigma_2)}$$

We put values of σ_1 and σ_3 where $\sigma_1 = \pm 1$ and $\sigma_3 = \pm 1$

$$\Rightarrow Z = \sum_{\sigma_2} e^{h'\sigma_2} \left(e^{1(h + v\bar{\sigma} + v\sigma_2) + 1(h + v\bar{\sigma} + v\sigma_2)} + e^{-1(h + v\bar{\sigma} + v\sigma_2) - 1(h + v\bar{\sigma} + v\sigma_2)} + e^{1(h + v\bar{\sigma} + v\sigma_2) - 1(h + v\bar{\sigma} + v\sigma_2)} + e^{-1(h + v\bar{\sigma} + v\sigma_2) + 1(h + v\bar{\sigma} + v\sigma_2)} \right)$$

$$\Rightarrow Z = \sum_{\sigma_2} e^{h'\sigma_2} \left(e^{2(h + v\bar{\sigma} + v\sigma_2)} + e^{-2(h + v\bar{\sigma} + v\sigma_2)} + 2 \right) \text{ then put } \sigma_2 = \pm 1$$

$$\Rightarrow Z = e^{h'} \left(e^{2(h + v\bar{\sigma} + v)} + e^{-2(h + v\bar{\sigma} + v)} + 2 \right) + e^{-h'} \left(e^{2(h + v\bar{\sigma} - v)} + e^{-2(h + v\bar{\sigma} - v)} + 2 \right)$$

$$\Rightarrow \boxed{Z = 4 \left[e^{h'} \cosh^2(h + v\bar{\sigma} + v) + e^{-h'} \cosh^2(h + v\bar{\sigma} - v) \right]} \text{ --- (2)}$$

This equation is the partition function Z of the system, to get $\bar{\sigma}$ we start the following steps:

$$\langle \sigma_1 + \sigma_3 \rangle = \frac{1}{Z} \frac{\partial Z}{\partial h} = \frac{\cancel{2} \left[e^{h'} \cosh(h + v\bar{\sigma} + v) \sinh(h + v\bar{\sigma} + v) + e^{-h'} \cosh(h + v\bar{\sigma} - v) \sinh(h + v\bar{\sigma} - v) \right]}{\cancel{4} \left[e^{h'} \cosh^2(h + v\bar{\sigma} + v) + e^{-h'} \cosh^2(h + v\bar{\sigma} - v) \right]} \quad \text{--- (3)}$$

Also

$$\langle \sigma_2 \rangle = \frac{1}{Z} \frac{\partial Z}{\partial h'} = \frac{\left[e^{h'} \cosh^2(h + v\bar{\sigma} + v) - e^{-h'} \cosh^2(h + v\bar{\sigma} - v) \right]}{\left[e^{h'} \cosh^2(h + v\bar{\sigma} + v) + e^{-h'} \cosh^2(h + v\bar{\sigma} - v) \right]} \quad \text{--- (4)}$$

but $\langle \sigma_1 + \sigma_3 \rangle = 2\langle \sigma_2 \rangle$ so from Eqs(3) and(4) we get

$$\begin{aligned} & \frac{\cancel{2} \left[e^{h'} \cosh(h + v\bar{\sigma} + v) \sinh(h + v\bar{\sigma} + v) + e^{-h'} \cosh(h + v\bar{\sigma} - v) \sinh(h + v\bar{\sigma} - v) \right]}{\left[e^{h'} \cosh^2(h + v\bar{\sigma} + v) + e^{-h'} \cosh^2(h + v\bar{\sigma} - v) \right]} \\ &= \cancel{2} \frac{\left[e^{h'} \cosh^2(h + v\bar{\sigma} + v) - e^{-h'} \cosh^2(h + v\bar{\sigma} - v) \right]}{\left[e^{h'} \cosh^2(h + v\bar{\sigma} + v) + e^{-h'} \cosh^2(h + v\bar{\sigma} - v) \right]} \end{aligned}$$

when $h = h' = 0$ we obtain

$$\left[\cosh(v\bar{\sigma} + v) \sinh(v\bar{\sigma} + v) + \cosh(v\bar{\sigma} - v) \sinh(v\bar{\sigma} - v) \right] = \left[\cosh^2(v\bar{\sigma} + v) - \cosh^2(v\bar{\sigma} - v) \right]$$

To solve this equation as polynomial equation we use these following substitutions

$x = e^{v\bar{\sigma}}$ and $y = e^{v\bar{\sigma}}$ in order to obtain

$$\frac{1}{4} \left[\frac{\cosh(v\bar{\sigma} + v) \sinh(v\bar{\sigma} + v)}{(xy + x^{-1}y^{-1})} + \frac{\cosh(v\bar{\sigma} - v) \sinh(v\bar{\sigma} - v)}{(xy^{-1} + x^{-1}y)} \right] = \frac{1}{4} \left[\frac{\cosh^2(v\bar{\sigma} + v)}{(xy + x^{-1}y^{-1})^2} - \frac{\cosh^2(v\bar{\sigma} - v)}{(xy^{-1} + x^{-1}y)^2} \right]$$

$$\Rightarrow (xy + x^{-1}y^{-1})(xy - x^{-1}y^{-1}) + (xy^{-1} + x^{-1}y)(xy^{-1} - x^{-1}y) = (xy + x^{-1}y^{-1})^2 - (xy^{-1} + x^{-1}y)^2$$

$$\Rightarrow (x^2y^2 - x^{-2}y^{-2}) + (x^2y^{-2} - x^{-2}y^2) = x^2y^2 + x^{-2}y^{-2} + 2 - x^2y^{-2} - x^{-2}y^2 - 2$$

$$\Rightarrow \cancel{x^2y^2} - x^{-2}y^{-2} + x^2y^{-2} - \cancel{x^{-2}y^2} = \cancel{x^2y^2} + x^{-2}y^{-2} + \cancel{2} - x^2y^{-2} - \cancel{x^{-2}y^2} - \cancel{2}$$

$$\Rightarrow -x^{-2}y^{-2} + x^2y^{-2} = +x^{-2}y^{-2} - x^2y^{-2}$$

$$-2x^{-2}y^{-2} + 2x^2y^{-2} = 0$$

$$(x^2 + x^{-2})y^{-2} = 0$$

The trivial solution is obtained by

$$(x^2 - x^{-2})y^{-2} = 0 \quad \Rightarrow (x^4 - 1)y^{-2} = 0$$

$x = 1$ only because $x = e^{\nu\bar{\sigma}} = 1$ this means that $\nu\bar{\sigma} = 0$ but $\nu \neq 0$ so $\bar{\sigma} = 0$

where $\nu \neq 0$ (ν is spin-spin interacting energy)

also $y \neq 0$ because $y = e^{\nu}$ and $\nu \neq 0$.

So the Bethe- peierls approximation does not predict a spontaneous magnetization for any value of ν .