Provided for non-commercial research and education use. Not for reproduction, distribution or commercial use.



This article appeared in a journal published by Elsevier. The attached copy is furnished to the author for internal non-commercial research and education use, including for instruction at the authors institution and sharing with colleagues.

Other uses, including reproduction and distribution, or selling or licensing copies, or posting to personal, institutional or third party websites are prohibited.

In most cases authors are permitted to post their version of the article (e.g. in Word or Tex form) to their personal website or institutional repository. Authors requiring further information regarding Elsevier's archiving and manuscript policies are encouraged to visit:

http://www.elsevier.com/copyright

Journal of Electron Spectroscopy and Related Phenomena 175 (2009) 21-26

Contents lists available at ScienceDirect

# Journal of Electron Spectroscopy and Related Phenomena

journal homepage: www.elsevier.com/locate/elspec



# XPS and magnetic studies of vanadium tellurite glassses

# A. Mekki<sup>a,\*</sup>, G.D. Khattak<sup>a</sup>, L.E. Wenger<sup>b</sup>

<sup>a</sup> Physics Department, King Fahd University of Petroleum & Minerals, Dhahran 31261, Saudi Arabia <sup>b</sup> Department of Physics, The University of Alabama at Birmingham, Birmingham, AL 35294-1170, USA

## ARTICLE INFO

Article history: Received 28 February 2009 Received in revised form 4 July 2009 Accepted 5 July 2009 Available online 18 August 2009

*Keywords:* Vanadium tellurite glasses XPS study of vanadate glasses Magnetic susceptibility of vanadate glasses

## ABSTRACT

A series of tellurite glasses containing vanadium with nominal composition  $xV_2O_5-(1-x)TeO_2$ , where x = 0.1, 0.2, 0.3, and 0.4, have been prepared by melt quenching technique and investigated using Xray photoelectron spectroscopy (XPS) and magnetization techniques. The Te 3d core level spectra show asymmetry and therefore were fitted with two contributions one from Te ions in TeO<sub>4</sub> trigonal bipyramid (tbp) configuration and the other from Te ions in TeO<sub>3</sub> trigonal pyramid (tp) configuration. Quantitative analysis of the areas of both components support the model of a change of  $TeO_4$  to  $TeO_3$  to upon  $V_2O_5$  addition. The V 2p core level spectra did not show any asymmetry and were fitted with a single component corresponding to V<sup>5+</sup> while the O 1s core level spectra show slight asymmetry on the higher binding energy side of the main peak and were fitted with two contributions, one due to bridging oxygen (BO) atoms and the other due to non-bridging oxygen (NBO) atoms. The fraction of NBO, determined from these spectra, is found to increase with increasing  $V_2O_5$  content suggesting the formation of TeO<sub>3</sub> units and also that V<sub>2</sub>O<sub>5</sub> is acting as a glass modifier in this glass series. The magnetic susceptibility data for these glasses follow a Curie–Weiss behavior  $(M/H = C/(T - \theta))$  which also indicates the presence of some V ions existing in a magnetic state, i.e., a valence state other than that of the non-magnetic state  $V^{5+}$ . The ratios of  $V^{4+}/V_{total}$  as determined from magnetic susceptibility measurements assuming only the presence of magnetic V<sup>4+</sup> and non-magnetic V<sup>5+</sup> ions vary from 0.030 to 0.037. These ratios are too small to be detected by XPS measurements and this is why no asymmetry is observed in the V2p spectra of these glasses.

© 2009 Elsevier B.V. All rights reserved.

## 1. Introduction

Tellurite glasses are high-index optical glasses that permit high levels of infrared radiation to be transmitted and thus have potential applicability as materials for acousto-optical devices and lasers, or as photochromic glasses. Tellurium dioxide, like vanadium pentoxide, is a conditional glass former and is the basis for forming the tellurite glasses. From various spectroscopic methods such as infrared [1–3], Raman [3–7], nuclear magnetic resonance [8], and X-ray absorption spectroscopy [9–11] as well as from X-ray [11,12] and neutron diffraction [13,14] techniques, it has been established that the basic structure of the tellurite glasses is a TeO<sub>4</sub> trigonal bipyramid (tbp) with a lone pair of electrons in one of its equatorial sites. Tellurium oxide under normal conditions does not have the ability to form a glass structure easily without a modifier like an alkali oxide, an alkaline-earth oxide, a transition-metal oxide, or another glass former [13,15]. Researchers have different opinions as to why it is so difficult to form a glass of pure TeO<sub>2</sub>. Neov et al. [13,15] hypothesized that TeO<sub>2</sub> could not form a glass by itself because the Te–O bond is too strongly covalent to permit the requisite amount of distortion for a glass structure. Another prevalent view [16,17] is that the repulsive forces due to the lone pair of electrons resist the free movement of the trigonal bipyramid in space during the cooling of the melt and hence the formation of glass. In a binary tellurite glass the effect of the lone pair of electrons is limited by the introduction of new structural units that are compatible with TeO<sub>4</sub> (tbp) and thus the glass formation becomes easier.

The addition of another conditional glass former such as V<sub>2</sub>O<sub>5</sub> to TeO<sub>2</sub> has been extensively studied by several authors. For example, Sakida et al. [18] studied the structure of V<sub>2</sub>O<sub>5</sub>–TeO<sub>2</sub> glasses by means of <sup>125</sup>Te and <sup>51</sup>V static NMR spectroscopies. They found that the fraction of TeO<sub>3</sub> and the NBO sites increase with increasing V<sub>2</sub>O<sub>5</sub> content. Hoppe et al. [19] used neutron diffraction to study the structure of this glass system. They found that in glasses of 10 mol% V<sub>2</sub>O<sub>5</sub> each VO<sub>n</sub> polyhedron is accompanied with a TeO<sub>3</sub> unit. The number of TeO<sub>3</sub> units increases further for glasses of higher V<sub>2</sub>O<sub>5</sub> fractions but it remains less than the number of V sites. The results are largely in agreement with those of Sakida et al. However, both authors did not estimate the fraction of various oxidation states in the glasses.

X-ray photoelectron spectroscopy (XPS) has been found to be a useful technique in assessing the local glass structure and even

<sup>\*</sup> Corresponding author. Tel.: +966 38604292; fax: +966 38602293. *E-mail address*: akmekki@kfupm.edu.sa (A. Mekki).

<sup>0368-2048/\$ -</sup> see front matter © 2009 Elsevier B.V. All rights reserved. doi:10.1016/j.elspec.2009.07.001

#### A. Mekki et al. / Journal of Electron Spectroscopy and Related Phenomena 175 (2009) 21-26

#### Table 1

#### Nominal and actual composition (molar fraction) of $xV_2O_5 - (1 - x)TeO_2$ .

Nominal		Actual (from ICI	2)
V <sub>2</sub> O <sub>5</sub>	TeO <sub>2</sub>	V <sub>2</sub> O <sub>5</sub>	TeO <sub>2</sub>
0.10	0.90	0.09	0.91
0.20	0.80	0.18	0.82
0.30	0.70	0.27	0.73
0.40	0.60	0.36	0.64

The relative uncertainty in the ICP results is  $\pm 5\%$ 

estimating the ratio of the different valence states in the TM-oxide glasses [20]. In this study, the state of V ions in vanadium-tellurite glasses will be presented from XPS spectra in combination with magnetization measurement data. We will also use XPS to measure the NBO and TeO<sub>3</sub> content in these glasses and compare our findings with those of previous work.

### 2. Experimental procedure

## 2.1. Glass preparation

All glasses were prepared by melting dry mixtures of reagent grade V<sub>2</sub>O<sub>5</sub> and TeO<sub>2</sub> in alumina crucibles to form batch  $xV_2O_5 - (1 - x)TeO_2$  compositions with x = 0.1, 0.2, 0.3, and 0.4. Since oxidation and reduction reactions in a glass melt are known to depend on the size of the melt, on the sample geometry, on whether the melt is static or stirred, on thermal history, and on quenching rate, all glass samples were prepared under similar conditions to minimize these factors. Approximately 30 g of chemicals were thoroughly mixed to obtain a homogenized mixture for each  $V_2O_5$  concentration. The crucible containing the batch mixture was then placed electrically heated melting furnace maintained at 850-900 °C. The melt was left for about 1 h under atmospheric conditions in the furnace during which the melt was occasionally stirred with an alumina rod. The homogenized melt was then cast onto a stainless steel plate mold to form glass buttons and glass rods of approximately 5-mm diameter for XPS measurements. Xray powder diffraction analysis indicated that the glasses formed were completely amorphous. After preparation, the samples were stored in a vacuum dessicator to minimize any further oxidation of the glass samples. The actual compositions of the glasses were determined by inductively coupled plasma spectroscopy (ICP) and are listed in Table 1. Although the inclusion of alumina from the crucibles used in the melting of the glass mixtures is a possible source of impurities, no signals for aluminum were detected in both the XPS and ICP studies of these glasses.

## 2.2. X-ray photoelectron spectroscopy (XPS)

High-resolution photoelectron spectra were collected on a VG scientific ESCALAB MKII spectrometer equipped with dual aluminium-magnesium anode using Al K( radiation (hv = 1486.6 eV). The electron analyzer was set at a pass energy of 20 eV. The energy scale of the spectrometer was calibrated using the core level of Cu  $2p_{3/2}\ (932.4\,eV)$  and the energy separation between Cu  $2p_{3/2}$  and Cu  $2p_{1/2}$  (19.8 eV). For self-consistency, the C 1s line of 284.6 eV binding energy was used as a reference for all charge shift corrections as this peak arises from hydrocarbon contamination and its binding energy is generally accepted as remaining constant, irrespective of the chemical state of the sample. This hydrocarbon contamination was further enhanced as the XPS measurements had to be performed on uncleaved surfaces of the glasses as the glass rod samples shattered under vacuum cleavage. The typical time required to collect an XPS spectrum for a sample was about 4h and the base pressure in

the analysis chamber during these measurements was less than  $10^{-10}$  mbar.

### 2.3. Magnetization measurements

Temperature-dependent dc magnetic susceptibilities were measured using a Quantum Design SQUID magnetometer (model MPMS-5S) in a magnetic field of 5000 Oe over a temperature range of 5–300 K. Each glass sample was suspended by thread in order to eliminate any sample holder correction to the measured magnetization. The overall accuracy of the magnetic measurements is estimated to be approximately 3% due to the uncertainty of the magnetometer calibration.

#### 3. Results

Fig. 1 shows high resolution Te 3d spin-orbit doublet spectra for the tellurite containing glasses. The peak intensity decreases with increasing  $V_2O_5$  content in the glasses, while the peak position for Te  $3d_{5/2}$  at ~576.5 eV does not vary much with composition. The peak width at half maximum (FWHM) shows first an increase, then a decrease with increasing V<sub>2</sub>O<sub>5</sub> content. This variation might indicate the change of TeO<sub>4</sub> tbp to TeO<sub>3</sub> tp configuration with the addition of V<sub>2</sub>O<sub>5</sub> content. The O 1s core level spectra are displayed in Fig. 2 for all compositions. A slight asymmetry is observed on the higher binding energy (BE) side of the main peak having an average value of  $\sim$  530.5 eV. This peak is also fitted with two contributions, one due to oxygen atoms in a bridging environment (BO) and the other due to oxygen atoms in a non-bridging environment (NBO). The O 1s spectra show an increase in intensity and this is well understood since we are substituting TeO<sub>2</sub> by V<sub>2</sub>O<sub>5</sub> and therefore introducing more oxygen atoms with increasing V2O5 content in the glass. Fig. 3 shows the V 2p spin-orbit doublet. The peak intensity of the V  $2p_{3/2}$  increases with increasing V<sub>2</sub>O<sub>5</sub> content in the glass but its BE does not vary much and has an average value of ~517.2 eV. The full-width at half maximum (FWHM) is also constant at  $\sim$ 1.9 eV for all glasses. This indicates that vanadium ions might exist mostly in a single oxidation state.

The temperature dependence of the magnetic susceptibility data for all glass samples appear to follow a Curie–Weiss behavior,  $M/H = C/(T - \theta)$ , with a temperature-independent background. After determining the background contribution from a high-temperature extrapolation of M/H-vs-1/T plots, the resulting  $M_{\text{corr}}/H = (M/H - (M/H)_{\text{background}})$  data were found to follow a Curie–Weiss behavior,  $M_{\text{corr}}/H = C/(T - \theta)$ , with the Curie constant



Fig. 1. Te 3d core level spectra for the vanadium tellurite glasses.

# Author's personal copy

#### A. Mekki et al. / Journal of Electron Spectroscopy and Related Phenomena 175 (2009) 21-26

Table 2

Magnetic susceptibility results for the copper telluride  $xV_2O_5-(1-x)TeO_2$  glasses and the concentration of V<sup>4+</sup> ions to total vanadium ions.

x (M/H) <sub>background</sub> (10 <sup>-7</sup> emu/Oe g)		$C(10^{-5} \mathrm{emu}\mathrm{K}/\mathrm{Oe}\mathrm{g})$	θ (K)	$[V^{4+}]/[V_{total}]$
0.10	-2.23	1.39	1.06	0.030
0.20	-2.05	3.08	-0.05	0.034
0.30	-0.71	5.22	1.24	0.039
0.40	-0.26	6.62	0.57	0.037

Table 3

Peak positions (in eV) for the core levels Te 3d, V 2p and O 1s relative to C 1s (284.6 eV), their corresponding full-width at half maximum (FWHM) and separation  $\Delta E$  of the spin orbit peaks  $xV_2O_5-(1-x)TeO_2$  glasses.

x	Te 3d <sub>5/2</sub> FWHM (eV)	$\Delta E$ Te 3d (eV)	V 2p <sub>3/2</sub> FWHM (eV)	$\Delta E \vee 2p (eV)$	O 1s FWHM (eV)
0.10	576.5	10.6	517.1	7.3	530.4
	2.45		2.0		2.5
0.20	576.5	10.4	517.1	7.3	530.5
	2.61		2.0		2.5
0.30	576.6	10.3	517.2	7.3	530.6
	2.11		1.8		1.9
0.40	576.7	10.4	517.3	7.3	530.6
	2.13		1.8		2.0
TeO <sub>2</sub>	576.1	10.4			530.3
	2.1				1.93
$V_2O_5$			517.3	7.3	530.2
			2.0		1.89

The uncertainty in the peak position is  $\pm 0.10 \text{ eV}$  and in FWHM is  $\pm 0.20 \text{ eV}$ .



Fig. 2. O 1s core level spectra for the vanadium tellurite glasses.



Fig. 3. V 2p core level spectra for the vanadium tellurite glasses.



**Fig. 4.** The inverse of the "corrected" magnetic susceptibility  $M_{\text{corr}}/H$  (= $M/H - M_{\text{background}}/H$ ) as a function of the temperature *T* for the vanadium tellurite glasses.

*C* and the paramagnetic Curie temperature  $\theta$  determined from the least-squares fits of the  $H/M_{corr}$ -vs-*T* data as shown in Fig. 4. Table 2 displays these parameters ((M/H)<sub>background</sub>, *C*, and  $\theta$ ) for each of the glass samples. From the Curie constant *C*, the relative fraction of V<sup>4+</sup> ( $\mu_{eff}$  = 1.8  $\mu_B$ ) was determined assuming that only V<sup>4+</sup> and V<sup>5+</sup> ( $\mu_{eff}$  = 0  $\mu_B$ ) ions were present in these glasses.

#### 4. Discussion

Although vanadium can exist in two oxidation states V<sup>4+</sup> and V<sup>5+</sup> in glasses but the observation that the peak position and the width of the V 2p<sub>3/2</sub> peak do not change much with the V<sub>2</sub>O<sub>5</sub> content in these glasses (see Table 3) suggests that vanadium ions probably exists in only one oxidation state. In a previous study on vanadium phosphate glasses [21] it was found that the core level spectrum of V<sup>5+</sup> has a 2p<sub>3/2</sub> BE of ~517.3 eV, while that of V<sup>4+</sup> is at a BE of ~516.1 eV. Thus fitting the V 2p<sub>3/2</sub> core level peak into two contributions, one from the V<sup>4+</sup> ions at a BE of ~516 eV and the other

#### A. Mekki et al. / Journal of Electron Spectroscopy and Related Phenomena 175 (2009) 21-26



**Fig. 5.** V  $2p_{3/2}$  core level spectra fitted with a single peak due to V<sup>5+</sup> for (a) x = 0.1 and (b) x = 0.4 glass compositions.

due to V<sup>5+</sup> ions at a BE of ~517.3 eV, should permit a more quantitative estimate of the relative V concentration for these two oxidation states of vanadium. The sum of two weighted Gaussian–Lorentzian peaks centered at approximately 516 and 517.3 eV is then fitted to the experimental data. It is clear from Fig. 5 that the V  $2p_{3/2}$  core level spectra for the x = 0.1 and 0.4 glasses can be satisfactorily fitted with a single Gaussian–Lorentzian peak centered at ~517.2 eV and having FHMW ~2.0 eV indicating that no V<sup>4+</sup> ions are present beyond the 5% detection limit of the technique. Similar quantitative fits are found for the other two glass compositions. The results of the fitting are summarized in Table 4.

As was mentioned earlier, it was found in previous studies that Te [19] ions exist in TeO<sub>4</sub> tbp and TeO<sub>3</sub> tp environments in vanadium tellurite glass samples with similar composition to the one studied here. We, therefore, tried to fit the Te  $3d_{5/2}$  core level spectra with two contributions, one due to Te ions in TeO<sub>4</sub> environment at a BE of ~577 eV with a FWHM of ~1.9 EV and the other due to

#### Table 4

Peak positions, FWHM, and relative concentration of V<sup>4+</sup> resulting from the curve fittings of the V 2p spectra for the  $xV_2O_5-(1-x)TeO_2$  glasses.

х	V 2p <sub>3/2</sub>		$[V^{4+}]/[V_{total}]^a$	
	Position (eV)	FWHM (eV)		
0.10	517.1	2.0	0	
0.20	517.1	2.0	0	
0.30	517.2	1.8	0	
0.40	517.3	1.8	0	

The experimental uncertainty is  $\pm 0.2 \text{ eV}$  in the energy measurement.

 $^{\rm a}\,$  If there are any  $V^{4+}$  ions, their content is beyond the detection limit of the XPS technique.

Te ions in TeO<sub>3</sub> environment at a BE of  $\sim$ 576 eV with a FWHM of  $\sim$ 1.9 eV. Fig. 6 shows such a fitting for the glass composition x = 0.3 and x = 0.4. Similar fitting was performed on the other two samples and the results are displayed in Table 5. We see from the table that



**Fig. 6.** Te 3d<sub>5/2</sub> core level spectra fitted with two contributions, one due to TeO<sub>4</sub> tbp configuration and the other due to TeO<sub>3</sub> tp configuration for (a) x = 0.3 and (b) x = 0.4 glass compositions.

# Author's personal copy

#### A. Mekki et al. / Journal of Electron Spectroscopy and Related Phenomena 175 (2009) 21-26

Та	ble	25

Peak positions, area, FWHM, peak separation, and relative concentration of TeO<sub>3</sub> resulting from the curve fittings of the Te  $3d_{5/2}$  spectra for the  $xV_2O_5-(1-x)TeO_2$  glasses.

x	Te 3d <sub>5/2</sub> (TeO <sub>3</sub> )	Te 3d <sub>5/2</sub> (TeO <sub>3</sub> )			Te 3d <sub>5/2</sub> (TeO <sub>4</sub> )			[TeO <sub>3</sub> ]%
	Position (eV)	FWHM (eV)	Area	Position (eV)	FWHM (eV)	Area		
0.10	575.7	1.86	4,524	576.7	2.00	18,751	1.0	19
0.20	575.7	1.90	12,799	576.7	1.85	29,030	1.0	31
0.30	576.1	1.75	8,060	577.1	1.85	11,381	1.0	41
0.40	576.3	1.80	31,712	577.3	1.80	28,798	1.0	52

The experimental uncertainty is  $\pm 0.2$  eV in the energy measurement and  $\pm 10\%$  in the relative concentration of TeO<sub>3</sub>.



Fig. 7. O 1s core level spectra fitted with BO and NBO peaks for (a) x = 0.1 and (b) x = 0.4 glass compositions.

substituting V<sub>2</sub>O<sub>5</sub> for TeO<sub>2</sub> increases the TeO<sub>3</sub> content in the glasses from ~20% at the lowest V<sub>2</sub>O<sub>5</sub> content to ~50% at the highest V<sub>2</sub>O<sub>5</sub> content. These values are in reasonably good agreement with those found using NMR [18] and neutron diffraction [19].

Among the most informative XPS spectra in glasses are those related to the oxygen atoms, namely of the O 1s spectra. The XPS technique has been successfully used to determine the proportions of bridging oxygen atoms (BO) and non-bridging oxygen atoms (NBO) in various glass systems. It has been used in simple binary glasses as well as ternary glass systems. Correspondingly, the O 1s peaks for these  $V_2O_5$ -TeO<sub>2</sub> glasses may arise from oxygen atoms existing in some or all of the following structural units: Te–O–Te, Te–O–V, V–O–V, and V=O. Oxygen atoms that are covalently bonded to glass former atoms on both sides are typically called bridging oxygen (e.g., P–O–P, Si–O–Si, Ge–O–Ge), while oxygen atoms ionically bonded, at least on one side, or double bonded to a glass former atom are referred to as non-bridging oxygens and have lower binding energies (e.g., Si–O–Na, Ge–O–Na, V=O, P=O). As described in the previous section, the O 1s peaks for the

V<sub>2</sub>O<sub>5</sub>-TeO<sub>2</sub> glass samples exhibit an asymmetry and correspondingly the O 1s spectra were fitted to two Gaussian–Lorentzian peaks as shown in Fig. 7. The resulting relative abundance of NBO/O<sub>total</sub> (see Table 6) increased from 43% to 76% as the nominal V<sub>2</sub>O<sub>5</sub> content increased from 0.1 to 0.4. However since the abundance of the V–O–V and V–O–Te structures should increase with increasing V<sub>2</sub>O<sub>5</sub> content, the oxygen atoms in V–O–V and V–O–Te structures in these glasses are more likely to be associated with non-bridging oxygens (NBO) rather than BO atoms in order to account for the observed increase in the relative abundance of NBO. Furthermore, each TeO<sub>3</sub> unit introduced in the glass structure contributes 2 NBOs [19]. The NBO content introduced from TeO<sub>4</sub> units through the O<sub>3/2</sub>Te–O<sup>-</sup> structure are taken from Table 5 of Ref. [18]. Assuming that all V atoms in V<sub>2</sub>O<sub>5</sub> give rise to an NBO [21], the total number of NBOs in each glass composition would be given by:

$$\frac{\text{NBO}}{\text{TO}} = \frac{5x + 2(1 - x)[\text{TeO}_3] + y[\text{TeO}_4]}{2 + 3x}$$
(1)

Table 6

Peak positions, area, FWHM, peak separation, and relative concentration of NBO resulting from the curve fitting of the O 1s core level for the  $xV_2O_5-(1-x)TeO_2$  glasses.

x	O 1s (eV)	0 1s (eV)		I	$\Delta E_{\rm BO-NBO}$ (eV)	[NBO] %	
	NBO (area)	BO (area)	NBO	BO		Measured	Cal <sup>a</sup>
0.10	530.0 (3147)	531.2 (4248)	1.90	2.00	1.2	43	40
0.20	530.0 (7451)	531.2 (7892)	2.00	2.00	1.2	49	60
0.30	530.3 (5292)	531.3 (2630)	1.85	1.85	1.0	67	80
0.40	530.5 (25,624)	531.5 (7976)	1.90	1.90	1.0	76	

The experimental uncertainty is  $\pm 0.2 \text{ eV}$  in the energy measurement and  $\pm 10\%$  in the relative concentration of NBO.

<sup>a</sup> Calculated using Eq. (1).

### A. Mekki et al. / Journal of Electron Spectroscopy and Related Phenomena 175 (2009) 21-26

where the actual values of x as determined from ICP (Table 1) were used in Eq. (1) and  $[TeO_3]$  and  $[TeO_4]$  are the concentrations of  $TeO_3$ and TeO<sub>4</sub> structural units in the glass determined from XPS (Table 5). Each V<sub>2</sub>O<sub>5</sub> molecule would contribute five NBO atoms to the O 1s peak, each TeO<sub>3</sub> structural unit would result in 2 NBOs while the contribution of TeO<sub>4</sub> to the NBOs is taken from Table 5 of Ref. [18]. The qualitative agreement between the values of NBO/O<sub>total</sub> determined from the relative areas under the O 1s peaks are displayed in Table 6 and the corresponding values calculated from Eq. (1) strongly supports this conjecture that the oxygen atoms in V-O-V and V-O-Te are probably more non-bridging in character compared to those in Te-O-Te units. Moreover, these rather large relative abundances of NBO indicate that V<sub>2</sub>O<sub>5</sub> plays the role of a network modifier rather than a glass former in this glass system in a similar way it did in the vanadium phosphate glasses investigated by the same group [21]. We can in general state that  $V_2O_5$  can be considered as both a glass former and modifier.

## 5. Conclusions

XPS and magnetization studies were performed to investigate the effect of substituting various amounts of V<sub>2</sub>O<sub>5</sub> for TeO<sub>2</sub> in a  $V_2O_5$ -TeO<sub>2</sub> binary glass system. The XPS analysis of the Te  $3d_{5/2}$ core level indicated the existence of Both TeO<sub>4</sub> and TeO<sub>3</sub> and that the fraction of TeO<sub>3</sub> increases with increasing V<sub>2</sub>O<sub>5</sub> content in agreement with previous studies conducted on the same glass system. The O1s spectra with a conventional two-peak fit and the number of NBO in each glass was found to correspond to contributions from V-O-V, V-O-Te, V=O and TeO<sub>3</sub> and some of the TeO<sub>4</sub> structural units. The V 2p spectra were fitted with a single peak corresponding to V<sup>5+</sup> ions. The magnetic susceptibility data for these glasses follow a Curie-Weiss behavior which also indicates the presence of some V ions existing in a magnetic state V<sup>4+</sup> in addition to the presence of the non-magnetic state  $V^{5+}$ . This ~3.5% presence of V ions in V<sup>4+</sup> state is beyond the 5% detection limit of XPS.

#### Acknowledgments

Two of the authors (AM) and (GDK) would like to acknowledge the support of the Physics Department at KFUPM.

#### References

- [1] N. Uchida, K. Takahashi, K. Nakata, S. Shibusawa, Yogyo Kyokai Shi 86 (1978) 317.
- [2] T. Yoko, K. Kamiya, K. Tanaka, H. Yamada, S. Sakka, Nippon Seramikkusu Kyokai Gakujutsu Ronbunshi 97 (1989) 289.
- [3] J. Heo, D. Lam, G.H. Sigel Jr., E.A. Mendoza, D.A. Hensley, J. Am. Ceram. Soc. 75 (1992)277T. Sekiya, N. Mchida, A. Ohtsuka, M. Tonokawa, Nippon Seramikkusu Kyokai
- [4] Gakujutsu Ronbunshi 97 (1989) 1435. [5] T. Sekiya, N. Mchida, A. Ohtsuka, M. Tonokawa, J. Non-Cryst. Solids 144 (1992)
- 128.
- [6] Y. Himei, A. Osaka, T. Nanba, Y. Miura, J. Non-Cryst. Solids 177 (1994) 164.
- [7] M. Tastumisago, T. Minami, Y. Kowada, H. Adachi, Phys. Chem. Glasses 35 (1994) 89
- [8] T. Yoko, M. Fujita, F. Miyaji, S. Sakka, Chem. Expr. 5 (1990) 549.
- [9] A. Osaka, Q. Jianrong, T. Nanba, Y. Miura, T. Yao, J. Non-Cryst. Solids 142 (1992) 81.
- [10] H. Yamamoto, H. Nasu, J. Matsuoka, K. Kamiya, J. Non-Cryst. Solids 170 (1994) 87
- [11] Y. Shimizugawa, T. Maeseto, S. Suehara, S. Inoue, A. Nukui, J. Mater. Res. 10 (1995) 405
- [12] G.W. Brady, J. Chem. Phys. 27 (1957) 300. [13]
- S. Neov, V. Kozhukharov, I. Gerasimova, K. Krezhov, B. Sidzhimov, J. Phys. C: Solid State Phys. 12 (1979) 2475
- [14] K. Suzuki, J. Non-Cryst. Solids 95-96 (1987) 15.
- [15] S. Neov, I. Gerasimova, K. Krezhov, B. Sidzhimov, V. Kozhukharov, Phys. Status Solidi A 47 (1978) 743.
- [16] Y. Dimitriev, V. Dimitrov, M. Arnaudov, J. Mater. Sci. 18 (1983) 1353.
- R.N. Sinclair, A.C. Wright, B. Bachra, Y.B. Dimitriev, V.V. Dimitrov, M.G. Arnaudov, [17] J. Non-Cryst. Solids 232-234 (1998) 38.
- [18] S. Sakida, S. Hayakawa, T. Yoko, J. Phys.: Condens. Matter 12 (2000) 2579. U. Hoppe, E. Yousef, C. Russel, J. Neuefeind, A.C. Hannon, Solid Sate Commun. [19]
- 123 (2002) 273
  - [20] A. Mekki, D. Holland, Kh.A. Ziq, C.F. McConville, J. Non-Cryst. Solids 272 (2000) 179.
  - [21] G.D. Khattak, A. Mekki, L.E. Wenger, J. Non-Cryst. Solids, in press.