

# Chapter 8

The time-dependent wave equation of a particle in a three dimensional box ( $L, L, L$ ) is

$$-\frac{\hbar^2}{2m} \nabla^2 \Psi(\vec{r}, t) + U(\vec{r}) \Psi(\vec{r}) = i\hbar \frac{\partial \Psi(\vec{r}, t)}{\partial t}$$

↑  
Laplacian;  $\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$

for stationary states  $\Psi(\vec{r}, t) = \Psi(\vec{r}) e^{-i\omega t}$

⇒ time-independent S.E. is

$$\left[ -\frac{\hbar^2}{2m} \nabla^2 + U(\vec{r}) \right] \Psi(\vec{r}) = E \Psi(\vec{r})$$

↑  
Hamiltonian operator

↑↑  
energy observable

$\Psi(\vec{r}) = \Psi(x) \cdot \Psi(y) \cdot \Psi(z)$  in Cartesian coordinates

In this case  $\Psi(x) = \sqrt{\frac{2}{L}} \sin\left(n_1 \frac{\pi}{L} x\right)$

$$\Psi(y) = \sqrt{\frac{2}{L}} \sin\left(n_2 \frac{\pi}{L} y\right)$$

$$\Psi(z) = \sqrt{\frac{2}{L}} \sin\left(n_3 \frac{\pi}{L} z\right)$$

So 
$$E = \frac{\pi^2 \hbar^2}{2mL^2} (n_1^2 + n_2^2 + n_3^2)$$

and 
$$\psi(x, y, z) = \left(\frac{2}{L}\right)^{3/2} \sin\left(\frac{n_1 \pi x}{L}\right) \sin\left(\frac{n_2 \pi y}{L}\right) \sin\left(\frac{n_3 \pi z}{L}\right)$$

with  $n_1 = 1, 2, \dots$

$n_2 = 1, 2, \dots$

$n_3 = 1, 2, \dots$

Degeneracy appears when different states have the same energy!

Example

$n_1$	$n_2$	$n_3$	
1	1	2	$\rightarrow \psi_{112} = \left(\frac{2}{L}\right)^{3/2} \sin\left(\frac{\pi x}{L}\right) \sin\left(\frac{\pi y}{L}\right) \sin\left(\frac{2\pi z}{L}\right)$
1	2	1	$\rightarrow \psi_{121} = \left(\frac{2}{L}\right)^{3/2} \sin\left(\frac{\pi x}{L}\right) \sin\left(\frac{2\pi y}{L}\right) \sin\left(\frac{\pi z}{L}\right)$
2	1	1	$\rightarrow \psi_{211} = \left(\frac{2}{L}\right)^{3/2} \sin\left(\frac{2\pi x}{L}\right) \sin\left(\frac{\pi y}{L}\right) \sin\left(\frac{\pi z}{L}\right)$

Three different state having the

same energy

$$E_{112} = E_{121} = E_{211} = \frac{6 \pi^2 \hbar^2}{2mL^2}$$

In the case of central forces acting on a particle we use spherical coordinates for the Laplacian.

Schrodinger equation is

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} + \frac{2\mu}{\hbar^2} [E - U] \psi = 0 \quad \text{--- (1)}$$

The wave function  $\psi(r, \theta, \phi) = R(r) \Theta(\theta) \Phi(\phi)$

this leads to three differential equations

$$\frac{d^2 \Phi}{d\phi^2} + m_l^2 \Phi = 0 \quad \text{--- (2)}$$

$$\frac{1}{\sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{d\Theta}{d\theta} \right) + \left[ l(l+1) - \frac{m_l^2}{\sin^2 \theta} \right] \Theta = 0 \quad \text{--- (3)}$$

$$\frac{1}{r^2} \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right) + \frac{2\mu}{\hbar^2} \left[ E - U - \frac{\hbar^2}{2m} \frac{l(l+1)}{r^2} \right] R = 0 \quad \text{--- (4)}$$

Solving the Schrodinger equation (1) means solving the above three equations.

Solution of equation (2) is  $\Phi(\phi) = e^{im_l \phi}$

Solution of equation (3) are the "associated Legendre polynomials". See Table 7.2 in your textbook. They are called  $P_l^{m_l}(\cos\theta)$

The product  $\Theta(\theta)\Phi(\phi)$  specify the full angular dependence of the central force wavefunction and are known as "spherical harmonics", denoted  $Y_l^{m_l}(\theta, \phi)$ . See Table 7.3 in your text book.

In our case  $|\vec{L}|$ ,  $L_z$  and  $E$  are sharp observables

$$|\vec{L}| = \sqrt{l(l+1)} \hbar \quad l = 0, 1, 2, \dots$$

$$L_z = m_l \hbar \quad -l \leq m_l \leq l$$

$l$ : "orbital quantum number"

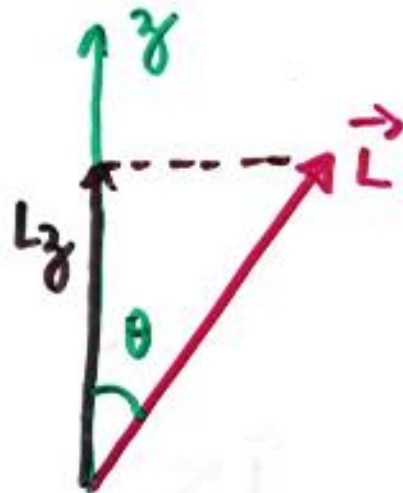
$m_l$ : "magnetic quantum number"

"Angular momentum and its  $z$ -component are QUANTIZED"

So:  $\Psi(\vec{r}) = R(r) Y_l^{m_l}(\theta, \phi)$

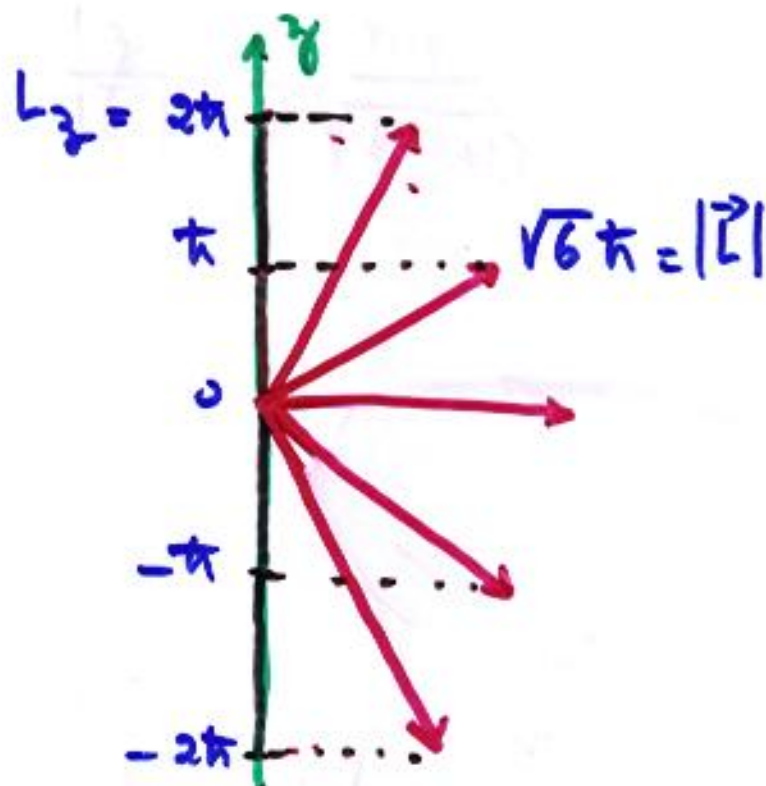
$\uparrow$  Total wavefunction       $\uparrow$  radial wave       $\uparrow$  angular wave

In quantum mechanics, space is "quantized".  
That is the angular momentum can take only  
specific directions with respect to the z-axis



$$\cos \theta = \frac{L_z}{|L|} = \frac{m_l \hbar}{\sqrt{l(l+1)} \hbar} \quad -l \leq m_l \leq l$$

For example:  $l = 2 \Rightarrow m_l = -2, -1, 0, 1, 2$ .



For hydrogen and hydrogen-like ions

The allowed energies are

$$E_n = - \frac{13.6 Z^2}{n^2} \quad n=1, 2, 3, \dots$$

$n$  is the "principle quantum number".

and the wave functions are

$$\psi_{n\ell m_\ell} = R_{n\ell}(r) Y_\ell^{m_\ell}(\theta, \phi)$$

↑  
given in Table 7.4 for  $n=1, 2$  and  $3$ .  
(your textbook)

For example

when  $n=1 \rightarrow \ell=0$  and  $m_\ell=0$

$$\psi_{100} = R_{10}(r) Y_0^0(\theta, \phi) \leftarrow \text{ground state wavefunction}$$

$$E_1 = -13.6 Z^2 \leftarrow \text{ground state energy}$$

From Tables  $\psi_{100} = \frac{1}{\sqrt{\pi}} \left(\frac{Z}{a_0}\right)^{3/2} e^{-\frac{Zr}{a_0}}$

For hydrogen atom:  $Z=1$

For helium  $\text{He}^+$  ion:  $Z=2$  etc...

$n$	shell symbol	$l$	subshell symbol
1	K	0	s
2	L	1	p
3	M	2	d
4	N	3	f
⋮	⋮	⋮	⋮

If  $n = 3$

- $l = 0$  ← 3s state
- $l = 1$  ← 3p state
- $l = 2$  ← 3d state

Always remember

$$n = 1, 2, 3, \dots$$

$$l = 0, 1, 2, \dots, (n-1)$$

$$-l \leq m_l \leq l$$

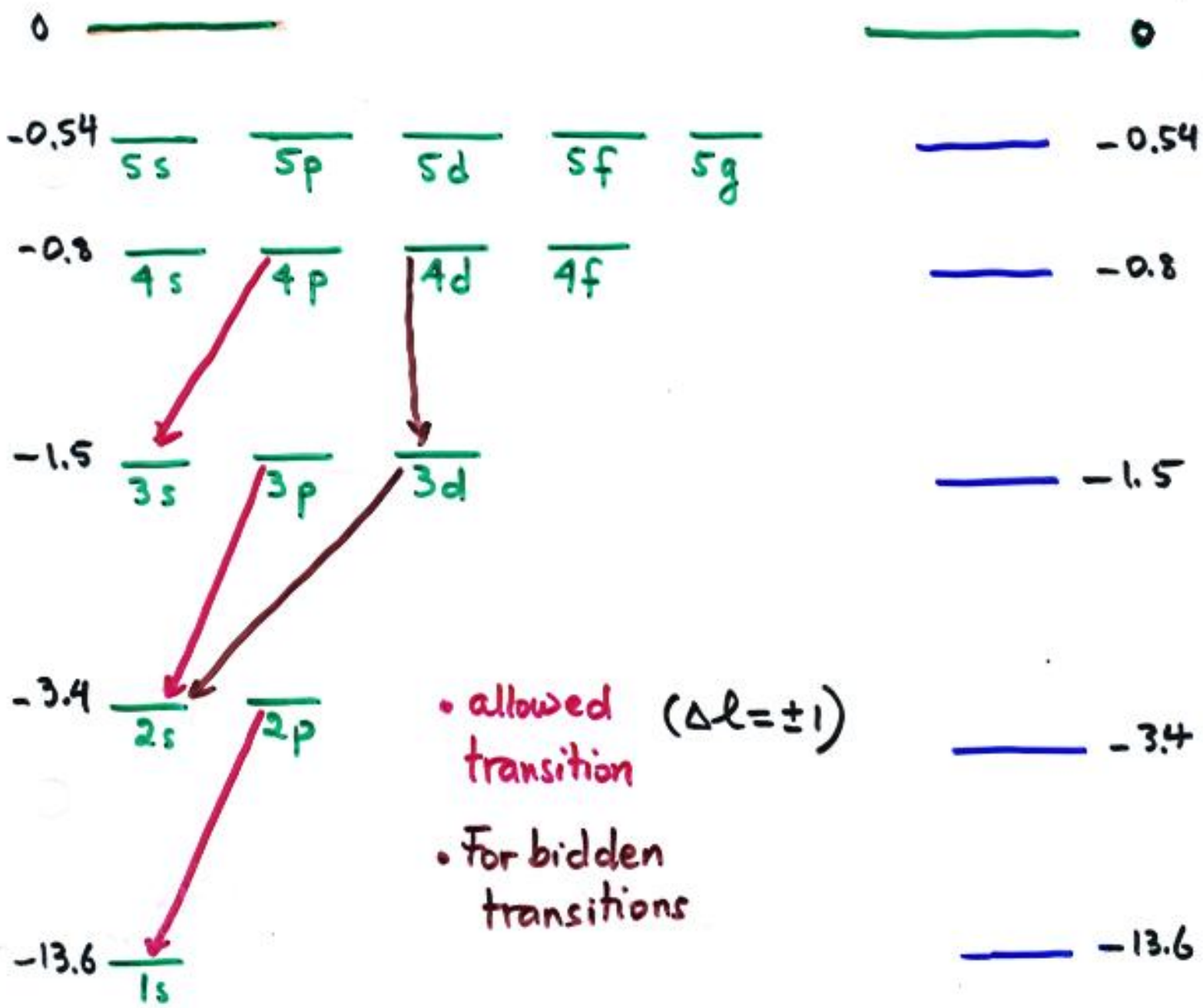
For  $n=2$   $l=0$   
 $l=1$

$m_l=0 \rightarrow \psi_{200}$   
 $m_l=-1 \rightarrow \psi_{21-1}$   
 $m_l=0 \rightarrow \psi_{210}$   
 $m_l=1 \rightarrow \psi_{211}$

$$E_2 = -\frac{13.4}{4} Z^2 = -3.4 Z^2 \text{ (eV)}$$

For hydrogen atoms  
 Q. M.

Bohr theory  $E(\text{eV})$



- allowed transition ( $\Delta l = \pm 1$ )
- For bidden transitions



- The ground state of the hydrogen and hydrogen like atoms will have quantum numbers:

$$n=1 \quad l=0 \quad m_l=0$$

$$E_1 = -13.6 Z^2 \text{ (eV)}$$

$$\psi_{100} = R_{10} Y_0^0 = \frac{1}{\sqrt{\pi}} \left( \frac{Z}{a_0} \right)^{3/2} e^{-\frac{Zr}{a_0}}$$

$\psi_{100}$  does not depend on  $\theta$  and  $\phi$  and is spherically symmetric so are all  $l=0$  (s-states).

The radial probability density for ANY state is

$$P(r) = r^2 |R(r)|^2$$

$R(r)$  is the radial wavefunction.

also  $\int_0^{\infty} P(r) dr = 1$  and  $\langle r \rangle = \int_0^{\infty} r P(r) dr$

↑  
average distance of the electron from the nucleus.

- The excited states of hydrogen like atoms are

$$n=2$$

$$l=0 \quad m_l=0 \quad E_2 = -3.4 Z^2 \text{ (eV)}$$

$$l=1 \quad m_l=1, 0, -1$$

$$\psi_{200} = \frac{1}{\sqrt{\pi}} \left(\frac{z}{2a_0}\right)^{3/2} \left(1 - \frac{zr}{2a_0}\right) e^{-\frac{zr}{2a_0}}$$

$$\psi_{211} = \frac{1}{\sqrt{\pi}} \left(\frac{z}{a_0}\right)^{3/2} \left(\frac{zr}{8a_0}\right) e^{-\frac{zr}{2a_0}} \sin\theta e^{i\phi}$$

$$\psi_{210} = \frac{1}{\sqrt{\pi}} \left(\frac{z}{2a_0}\right)^{3/2} \left(\frac{zr}{2a_0}\right) e^{-\frac{zr}{2a_0}} \cos\theta$$

$$\psi_{21-1} = -\frac{1}{\sqrt{\pi}} \left(\frac{z}{a_0}\right)^{3/2} \left(\frac{zr}{8a_0}\right) e^{-\frac{zr}{2a_0}} \sin\theta e^{-i\phi}$$

The level  $E_2$  is 4-fold degenerate.