

Global Existence and Nonexistence in a System of Petrovsky

Salim A. Messaoudi

*Mathematical Sciences Department, King Fahd University of Petroleum and Minerals,
Dhahran 31261, Saudi Arabia*

E-mail: messaoud@kfupm.edu.sa

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In this paper we consider the nonlinearly damped semilinear Petrovsky equation

$$u_{tt} + \Delta^2 u + au_t|u_t|^{m-2} = bu|u|^{p-2}$$

in a bounded domain, where $a, b > 0$. We prove the existence of a local weak solution and show that this solution blows up in finite time if $p > m$ and the energy is negative. We also show that the solution is global if $m \geq p$. © 2002 Elsevier Science

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1. INTRODUCTION

In [4], Guesmia considered the problem

$$\begin{aligned} u_{tt}(x, t) + \Delta^2 u(x, t) + q(x)u(x, t) + g(u_t(x, t)) &= 0, \quad x \in \Omega, \quad t > 0 \\ u(x, t) = \partial_\nu u(x, t) &= 0, \quad x \in \partial\Omega, \quad t \geq 0 \\ u(x, 0) = u_0(x), \quad u_t(x, 0) &= u_1(x), \quad x \in \Omega, \end{aligned} \quad (1.1)$$

where Ω is a bounded domain of $\mathbb{R}^n (n \geq 1)$, with a smooth boundary $\partial\Omega$, and ν is the unit outer normal on $\partial\Omega$. For g continuous, increasing, satisfying $g(0) = 0$, and $q: \Omega \rightarrow \mathbb{R}^+$, a bounded function, Guesmia [4] proved a global existence and a regularity result. He also established, under suitable growth conditions on g , decay results for weak, as well as strong, solutions. Precisely, the author showed that the solution decays exponentially if g behaves like a linear function, whereas the decay is of a polynomial order

otherwise. Results similar to the above system, coupled with a semilinear wave equation, have been established by Guesmia [5]. Also the system composed of the equation (1.1), with $\Delta^2 u_t(x, t) + \Delta g(\Delta u(x, t))$ in the place of $q(x)u(x, t) + g(u_t(x, t))$, has been treated by Aassila and Guesmia [1], and an exponential decay theorem, through the use of an important lemma of Komornik [6], has been established.

In this paper we are concerned with the problem

$$\begin{aligned} u_{tt} + \Delta^2 u + au_t|u_t|^{m-2} &= bu|u|^{p-2}, \quad x \in \Omega, \quad t > 0 \\ u(x, t) = \partial_\nu u(x, t) &= 0, \quad x \in \partial\Omega, \quad t \geq 0 \\ u(x, 0) = \phi(x), \quad u_t(x, 0) &= \varphi(x), \quad x \in \Omega, \end{aligned} \tag{1.2}$$

where $a, b > 0$ and $p, m > 2$. This is a problem similar to (1.1), which contains a nonlinear source term competing with the damping factor. We will establish an existence result and show that the solution continues to exist globally if $m \geq p$; however, it blows up in finite time if $m < p$. It is worth mentioning that it is only for simplicity that q is taken to be zero, $g(u_t(x, t)) = au_t|u_t|^{m-2}$, and the source term has a power form. The same theorems could be established for more general functions.

2. LOCAL EXISTENCE

In this section, we establish a local existence result for (1.2) under suitable conditions on m and p . First we consider, for v given, the linear problem

$$\begin{aligned} u_{tt} + \Delta^2 u + au_t|u_t|^{m-2} &= b|v|^{p-2}v, \quad x \in \Omega, \quad t > 0 \\ u(x, t) = \partial_\nu u(x, t) &= 0, \quad x \in \partial\Omega, \quad t > 0 \\ u(x, 0) = \phi(x), \quad u_t(x, 0) &= \varphi(x), \quad x \in \Omega, \end{aligned} \tag{2.1}$$

where u is the sought solution.

LEMMA 2.1. *Assume that*

$$\begin{aligned} 2 < p, \quad n &\leq 4 \\ 2 < p \leq 2(n - 2)/n - 4, \quad n &\geq 5. \end{aligned} \tag{2.2}$$

Then given any v in $C([0, T]; C_0^\infty(\Omega))$ and ϕ, φ in $C_0^\infty(\Omega)$, the problem (2.1) has a unique solution u satisfying

$$\begin{aligned} u \in L^\infty((0, T); W), \quad u_{tt} \in L^\infty((0, T); L^2(\Omega)) \\ u_t \in L^\infty((0, T); H_0^2(\Omega)) \cap L^m(\Omega \times (0, T)). \end{aligned} \tag{2.3}$$

Here $H_0^2(\Omega) = \{w \in H^2(\Omega) : w = \partial_\nu w = 0 \text{ on } \partial\Omega\}$ and $\mathbf{W} = \{w \in H^4(\Omega) \cap H_0^2(\Omega) : \Delta w = \partial_\nu \Delta w = 0 \text{ on } \partial\Omega\}$.

This lemma is a direct result of [7, Theorem 3.1, Chap. 1] (see also [2] and [4, Theorem 1.2]).

LEMMA 2.2. *Assume that (2.2) holds. Assume further that*

$$m \leq 2n/(n-4), \quad n \geq 5. \quad (2.4)$$

Then given any ϕ in $H_0^2(\Omega)$, φ in $L^2(\Omega)$, and v in $C([0, T]; H_0^2(\Omega))$, the problem (2.1) has a unique weak solution,

$$\begin{aligned} u &\in C([0, T]; H_0^2(\Omega)) \\ u_t &\in C([0, T]; L^2(\Omega)) \cap L^m(\Omega \times (0, T)). \end{aligned} \quad (2.5)$$

Moreover, we have

$$\begin{aligned} &\frac{1}{2} \int_{\Omega} [u_t^2 + (\Delta u)^2](x, t) dx + a \int_0^t \int_{\Omega} |u_t(x, s)|^m dx ds \\ &= \frac{1}{2} \int_{\Omega} [\varphi^2 + (\Delta \phi)^2](x) dx + b \int_0^t \int_{\Omega} |v|^{p-2} v u_t(x, s) dx ds, \\ &\quad \forall t \in [0, T]. \end{aligned} \quad (2.6)$$

Proof. We approximate ϕ , φ by sequences (ϕ^μ) , (φ^μ) in $C_0^\infty(\Omega)$, and v by a sequence (v^μ) in $C([0, T]; C_0^\infty(\Omega))$. We then consider the set of linear problems

$$\begin{aligned} u_{tt}^\mu + \Delta^2 u^\mu + a u_t^\mu |u_t^\mu|^{m-2} &= b |v^\mu|^{p-2} v^\mu, \quad x \in \Omega, \quad t > 0 \\ u^\mu(x, t) &= \partial_\nu u^\mu(x, t) = 0, \quad x \in \partial\Omega, \quad t > 0 \\ u^\mu(x, 0) &= \phi^\mu(x), \quad u_t^\mu(x, 0) = \varphi^\mu(x), \quad x \in \Omega. \end{aligned} \quad (2.7)$$

Lemma 2.1 guarantees the existence of a sequence of unique solutions (u^μ) satisfying (2.3). Now we proceed to show that the sequence (u^μ, u_t^μ) is Cauchy in

$$\mathbf{Y} := \{w : w \in C([0, T]; H_0^2(\Omega)), w_t \in C([0, T]; L^2(\Omega)) \cap L^m(\Omega \times (0, T))\}.$$

For this aim, we set

$$U := u^\mu - u^\nu, \quad V := v^\mu - v^\nu.$$

It is straightforward to see that U satisfies

$$\begin{aligned} U_{tt} + \Delta^2 U + a(u_t^\mu |u_t^\mu|^{m-2} - u_t^\nu |u_t^\nu|^{m-2}) &= b(|v^\mu|^{p-2} v^\mu - |v^\nu|^{p-2} v^\nu) \\ U(x, t) &= 0, \quad x \in \partial\Omega, \quad t > 0 \end{aligned} \quad (2.8)$$

$$U(x, 0) = U_0(x) = \phi^\mu(x) - \phi^\nu(x), \quad U_t(x, 0) = U_1(x) = \varphi^\mu(x) - \varphi^\nu(x).$$

We multiply Eq. (2.8) by U_t and integrate over $\Omega \times (0, t)$ to get

$$\begin{aligned} & \frac{1}{2} \int_{\Omega} [U_t^2 + (\Delta U)^2](x, t) dx + a \int_0^t \int_{\Omega} (u_t^\mu |u_t^\mu|^{m-2} - u_t^\nu |u_t^\nu|^{m-2}) U_t(x, s) dx ds \\ &= \frac{1}{2} \int_{\Omega} [U_1^2 + (\Delta U_0)^2](x) dx + b \int_0^t \int_{\Omega} [|v^\mu|^{p-2} v^\mu - |v^\nu|^{p-2} v^\nu] \\ & \quad \times U_t(x, s) dx ds. \end{aligned} \tag{2.9}$$

We then estimate the last term in (2.9) as follows:

$$\begin{aligned} & \int_{\Omega} \left| [|v^\mu|^{p-2} v^\mu - |v^\nu|^{p-2} v^\nu] U_t(x, s) \right| dx \\ & \leq C \|U_t\|_2 \|V\|_{2n/(n-4)} \left[\|v^\mu\|_{n(p-2)/2}^{p-2} + \|v^\nu\|_{n(p-2)/2}^{p-2} \right]. \end{aligned} \tag{2.10}$$

The Sobolev embedding and condition (2.2) give

$$\begin{aligned} & \|V\|_{2n/(n-2)} \leq C \|\Delta V\|_2, \\ & \|v^\mu\|_{n(p-2)/2}^{p-2} + \|v^\nu\|_{n(p-2)/2}^{p-2} \leq C [\|\Delta v^\mu\|_2^{p-2} + \|\Delta v^\nu\|_2^{p-2}], \end{aligned}$$

where C is a constant depending on Ω only. Therefore (2.10) takes the form

$$\begin{aligned} & \int_{\Omega} \left| [|v^\mu|^{p-2} v^\mu - |v^\nu|^{p-2} v^\nu] U_t(x, s) \right| dx \\ & \leq C \|U_t\|_2 \|\Delta V\|_2 [\|\Delta v^\mu\|_2^{p-2} + \|\Delta v^\nu\|_2^{p-2}]. \end{aligned}$$

Since $(u_t^\mu |u_t^\mu|^{m-2} - u_t^\nu |u_t^\nu|^{m-2})(u_t^\mu - u_t^\nu) \geq 0$ then (2.9) yields

$$\begin{aligned} & \frac{1}{2} \int_{\Omega} [U_t^2 + (\Delta U)^2](x, t) dx \leq \int_{\Omega} [U_1^2 + (\Delta U_0)^2](x) dx \\ & \quad + \Gamma \int_0^t \|U_t(\cdot, s)\|_2 \|\Delta V(\cdot, s)\|_2 ds, \end{aligned}$$

where Γ is a generic positive constant depending on C and the radius of the ball in $C([0, T]; H_0^2(\Omega))$ containing v^μ and v^ν . Young's inequality then gives

$$\begin{aligned} & \max_{0 \leq t \leq T} \int_{\Omega} [U_t^2 + (\Delta U)^2](x, t) dx \leq \Gamma \int_{\Omega} [U_1^2 + |\Delta U_0|^2](x) dx \\ & \quad + \Gamma T \max_{0 \leq t \leq T} \int_{\Omega} [V_t^2 + (\Delta V)^2](x, t) dx. \end{aligned}$$

Since (ϕ^μ) is Cauchy in $H_0^2(\Omega)$, (φ^μ) is Cauchy in $L^2(\Omega)$, and (v^μ) is Cauchy in $C([0, T]; H_0^2(\Omega))$, we conclude that (u^μ, u_t^μ) is Cauchy

in $C([0, T]; H_0^2(\Omega)) \times C([0, T]; L^2(\Omega))$. To show that u_t is Cauchy in $L^m(\Omega \times (0, T))$, we use

$$\|U_t\|_{L^m(\Omega \times (0, T))}^m \leq C \int_0^t \int_{\Omega} (u_t^\mu |u_t^\mu|^{m-2} - u_t^\nu |u_t^\nu|^{m-2}) U_t(x, s) dx ds, \quad (2.11)$$

which yields, by (2.9),

$$\begin{aligned} \|U_t\|_{L^m(\Omega \times (0, T))}^m &\leq \Gamma \int_{\Omega} [U_1^2 + (\Delta U_0)^2](x) dx \\ &\quad + \Gamma \int_0^T \|U_t(\cdot, s)\|_2 \|\Delta V(\cdot, s)\|_2 ds. \end{aligned}$$

Therefore (u_t^μ) is Cauchy in $L^m(\Omega \times (0, T))$ and hence (u^μ) is Cauchy in \mathbf{Y} . We now show that the limit u is a weak solution of (2.1) in the sense of [7]. That is for each θ in $H_0^2(\Omega)$ we must show that

$$\begin{aligned} \frac{d}{dt} \int_{\Omega} u_t(x, t) \theta(x) dx + \int_{\Omega} \Delta u(x, t) \Delta \theta(x) dx \\ + a \int_{\Omega} u_t |u_t|^{m-2}(x, t) \theta(x) dx = b \int_{\Omega} |v|^{p-2} v(x, t) \theta(x) dx, \end{aligned} \quad (2.12)$$

for almost all t in $[0, T]$. To establish this, we multiply Eq. (2.7) by θ and integrate over Ω , so we obtain

$$\begin{aligned} \frac{d}{dt} \int_{\Omega} u_t^\mu(x, t) \theta(x) dx + \int_{\Omega} \Delta u^\mu(x, t) \Delta \theta(x) dx \\ + a \int_{\Omega} u_t^\mu |u_t^\mu|^{m-2}(x, t) \theta(x) dx = b \int_{\Omega} |v^\mu|^{p-2} v^\mu(x, t) \theta(x) dx. \end{aligned} \quad (2.13)$$

As $\mu \rightarrow \infty$, we see that

$$\begin{aligned} \int_{\Omega} \Delta u^\mu(x, t) \Delta \theta(x) dx &\rightarrow \int_{\Omega} \Delta u(x, t) \Delta \theta(x) dx, \\ \int_{\Omega} |v^\mu|^{p-2} v^\mu(x, t) \theta(x) dx &\rightarrow \int_{\Omega} |v|^{p-2} v(x, t) \theta(x) dx \end{aligned}$$

in $C([0, T])$ and

$$\int_{\Omega} u_t^\mu |u_t^\mu|^{m-2}(x, t) \theta(x) dx \rightarrow \int_{\Omega} u_t |u_t|^{m-2}(x, t) \theta(x) dx$$

in $L^1((0, T))$. We thus have $\int_{\Omega} u_t(x, t) \theta(x) dx \{= \lim \int_{\Omega} u_t^\mu(x, t) \theta(x) dx\}$ is an absolutely continuous function on $[0, T]$, so (2.12) holds for almost all t in $[0, T]$. For the energy equality (2.6), we start from the energy equality for u^μ and proceed in the same way to establish it for u . To prove uniqueness,

we take v^μ and v^ν and let u^μ and u^ν be the corresponding solutions of (2.1). It is clear that $U = u^\mu - u^\nu$ satisfies

$$\begin{aligned} & \frac{1}{2} \int_{\Omega} [U_t^2 + (\Delta U)^2](x, t) dx + a \int_0^t \int_{\Omega} (u_t^\mu |u_t^\mu|^{m-2} - u_t^\nu |u_t^\nu|^{m-2}) U_t(x, s) dx ds \\ & = b \int_0^t \int_{\Omega} [|v^\mu|^{p-2} v^\mu - |v^\nu|^{p-2} v^\nu] U_t(x, s) dx ds. \end{aligned} \tag{2.14}$$

If $v^\mu = v^\nu$ then (2.14) shows that $U = 0$, which implies uniqueness. This completes the proof.

Remark 2.1. Note that the condition (2.4) on m is needed so that $\int_{\Omega} u_t^\mu |u_t^\mu|^{m-2}(x, t)\theta(x) dx$ and $\int_{\Omega} u_t |u_t|^{m-2}(x, t)\theta(x) dx$ make sense.

THEOREM 2.3. *Assume that (2.2) and (2.4) hold. Then given any ϕ in $H_0^2(\Omega)$, and φ in $L^2(\Omega)$, the problem (1.2) has a unique weak solution $u \in \mathbf{Y}$, for T small enough.*

Proof. For $M > 0$ large and $T > 0$, we define a class of functions $Z(M, T)$ which consists of all functions w in \mathbf{Y} satisfying the initial conditions of (1.2) and

$$\max_{0 \leq t \leq T} \frac{1}{2} \int_{\Omega} [w_t^2 + (\Delta w)^2](x, t) dx + a \int_0^T \int_{\Omega} |w_t(x, s)|^m dx ds \leq M^2. \tag{2.15}$$

$Z(M, T)$ is nonempty if M is large enough. This follows from the trace theorem (see [8]). We also define the map f from $Z(M, T)$ into \mathbf{Y} by $u := f(v)$, where u is the unique solution of the linear problem (2.1). We would like to show, for M sufficiently large and T sufficiently small, that f is a contraction from $Z(M, T)$ into itself.

By using the energy equality (2.5) we get

$$\begin{aligned} & \int_{\Omega} [u_t^2 + (\Delta u)^2](x, t) dx + 2a \int_0^t \int_{\Omega} |u_t(x, s)|^m dx ds \\ & \leq \int_{\Omega} [u_1^2 + (\Delta u_0)^2](x) dx + 2b \int_0^t \int_{\Omega} |v|^{p-1} |u_t|(x, s) dx ds, \quad \forall t \in [0, T] \\ & \leq \int_{\Omega} [u_1^2 + (\Delta u_0)^2](x) dx + 2b \int_0^t \|u_t\|_2 \|\Delta v\|_2^{p-1}, \quad \forall t \in [0, T]; \end{aligned} \tag{2.16}$$

consequently

$$\|u\|_{\mathbf{Y}}^2 \leq C \int_{\Omega} [u_1^2 + (\Delta u_0)^2](x) dx + CM^{p-1} T \|u\|_{\mathbf{Y}},$$

where C is independent of M . By choosing M large enough and T sufficiently small, (2.15) is satisfied; hence $u \in Z(M, T)$. This shows that f maps $Z(M, T)$ into itself.

Next we verify that f is a contraction. For this aim we set $U = u - \bar{u}$ and $V = v - \bar{v}$, where $u = f(v)$ and $\bar{u} = f(\bar{v})$. It is straightforward to see that U satisfies

$$\begin{aligned} U_{tt} + \Delta^2 U + a|u_t|^{m-2}u_t - a|\bar{u}_t|^{m-2}\bar{u}_t &= b|v|^{p-2}v - b|\bar{v}|^{p-2}\bar{v} \\ U(x, t) &= 0, \quad x \in \partial\Omega, \quad t > 0 \\ U(x, 0) &= U_t(x, 0) = 0, \quad x \in \Omega. \end{aligned} \quad (2.17)$$

By multiplying Eq. (2.17) by U_t and integrating over $\Omega \times (0, t)$, we arrive at

$$\begin{aligned} \int_{\Omega} [U_t^2 + (\Delta U)^2](x, t) dx + \int_0^t \int_{\Omega} (|u_t|^{m-2}u_t - |\bar{u}_t|^{m-2}\bar{u}_t)U_t(x, s) dx ds \\ \leq C \int_0^t \int_{\Omega} \|v\|^{p-2}v - \|\bar{v}\|^{p-2}\bar{v} \|U_t\|(x, s) dx ds. \end{aligned} \quad (2.18)$$

By using (2.2), (2.10), and (2.11), we obtain

$$\begin{aligned} \int_{\Omega} [U_t^2 + (\Delta U)^2](x, t) dx + \int_0^t \int_{\Omega} |U_t(x, s)|^m dx ds \\ \leq \Gamma \int_0^t \|U_t\|_2 \|\Delta V\|_2 (\|\Delta v\|_2^{p-2} + \|\Delta \bar{v}\|_2^{p-2})(\cdot, s) ds. \end{aligned}$$

Thus we have

$$\|U\|_{\mathbb{Y}}^2 \leq CTM^{p-2} \|V\|_{\mathbb{Y}}^2. \quad (2.19)$$

By choosing T so small that $\Gamma TM^{p-2} < 1$, (2.19) shows that f is a contraction. The contraction mapping theorem then guarantees the existence of a unique u satisfying $u = f(u)$. Obviously it is a solution of (1.2). The uniqueness of this solution follows from the energy inequality (2.18). The proof is completed.

3. BLOW-UP RESULT

In this section we show that the solution (2.5) blows up in finite time if $p > m$ and $E(0) < 0$, where

$$E(t) := \frac{1}{2} \int_{\Omega} [u_t^2 + (\Delta u)^2](x, t) dx - \frac{b}{p} \int_{\Omega} |u(x, t)|^p dx. \quad (3.1)$$

LEMMA 3.1. *Suppose that (2.2) holds. Then there exists a positive constant $C > 1$, depending on Ω only, such that*

$$\|u\|_p^s \leq C (\|\Delta u\|_2^2 + \|u\|_p^p) \quad (3.2)$$

for any $u \in H_0^2(\Omega)$ and $2 \leq s \leq p$.

Proof. If $\|u\|_p \leq 1$ then $\|u\|_p^s \leq \|u\|_p^2 \leq C\|\Delta u\|_2^2$ by Sobolev embedding theorems and the boundary conditions. If $\|u\|_p > 1$ then $\|u\|_p^s \leq \|u\|_p^p$. Therefore (3.2) follows.

We set

$$H(t) := -E(t) \tag{3.3}$$

and use, throughout this section, C to denote a generic positive constant depending on Ω only. As a result of (3.1)–(3.3), we have

COROLLARY 3.2. *Let the assumptions of the lemma hold. Then we have*

$$\|u\|_p^s \leq C(|H(t)| + \|u_t\|_2^2 + \|u\|_p^p) \tag{3.4}$$

for any $u \in H_0^2(\Omega)$ and $2 \leq s \leq p$.

THEOREM 3.3. *Let the conditions of the Theorem 2.3 be fulfilled. Assume further that*

$$E(0) < 0. \tag{3.5}$$

Then the solution (2.5) blows up in finite time.

Proof. We multiply Eq. (1.2) by $-u_t$ and integrate over Ω to get

$$H'(t) = a \int_{\Omega} |u_t(x, t)|^m dx \geq 0,$$

for almost every t in $[0, T)$ since $H(t)$ is absolutely continuous (see [2]); hence

$$0 < H(0) \leq H(t) \leq \frac{b}{p} \|u\|_p^p, \tag{3.6}$$

for every t in $[0, T)$, by virtue of (3.1) and (3.3). We then define

$$L(t) := H^{1-\alpha}(t) + \varepsilon \int_{\Omega} uu_t(x, t) dx \tag{3.7}$$

for ε small to be chosen later and

$$0 < \alpha \leq \min \left\{ \frac{(p-2)}{2p}, \frac{(p-m)}{p(m-1)} \right\}. \tag{3.8}$$

By taking a derivative of (3.7) and using Eq. (1.2) we obtain

$$\begin{aligned} L'(t) := & (1-\alpha)H^{-\alpha}(t)H'(t) + \varepsilon \int_{\Omega} [u_t^2 - (\Delta u)^2](x, t) dx \\ & + \varepsilon b \int_{\Omega} |u(x, t)|^p dx - a\varepsilon \int_{\Omega} |u_t|^{m-2} u_t u(x, t) dx. \end{aligned} \tag{3.9}$$

We then exploit Young's inequality,

$$XY \leq \frac{\delta^r}{r} X^r + \frac{\delta^{-q}}{q} Y^q, \quad X, Y \geq 0, \quad \delta > 0, \quad \frac{1}{r} + \frac{1}{q} = 1,$$

for $r = m$ and $q = m/(m-1)$ to estimate the last term in (3.9) as

$$\int_{\Omega} |u_t|^{m-1} |u| dx \leq \frac{\delta^m}{m} \|u\|_m^m + \frac{m-1}{m} \delta^{-m/(m-1)} \|u_t\|_m^m,$$

which yields, by substitution in (3.9),

$$\begin{aligned} L'(t) \geq & \left[(1-\alpha)H^{-\alpha}(t) - \frac{m-1}{m} \varepsilon \delta^{-m/(m-1)} \right] H'(t) \\ & + \varepsilon \int_{\Omega} [u_t^2 - (\Delta u)^2](x, t) dx + \varepsilon \left[pH(t) + \frac{P}{2} \int_{\Omega} [u_t^2 + (\Delta u)^2](x, t) dx \right] \\ & - \varepsilon a \frac{\delta^m}{m} \|u\|_m^m, \quad \forall \delta > 0. \end{aligned} \quad (3.10)$$

Of course (3.10) remains valid even if δ is time dependent, since the integral is taken over the x variable. Therefore by taking δ so that $\delta^{-m/(m-1)} = kH^{-\alpha}(t)$, for large k to be specified later, and substituting in (3.10) we arrive at

$$\begin{aligned} L'(t) \geq & \left[(1-\alpha) - \frac{m-1}{m} \varepsilon k \right] H^{-\alpha}(t) H'(t) \\ & + \varepsilon \left(\frac{P}{2} + 1 \right) \int_{\Omega} u_t^2(x, t) dx + \varepsilon \left(\frac{P}{2} - 1 \right) \int_{\Omega} (\Delta u(x, t))^2 dx \\ & + \varepsilon \left[pH(t) - \frac{k^{1-m}}{m} a H^{\alpha(m-1)}(t) \|u\|_m^m \right]. \end{aligned} \quad (3.11)$$

By exploiting (3.6) and the inequality $\|u\|_m^m \leq C \|u\|_p^m$, we obtain

$$H^{\alpha(m-1)}(t) \|u\|_m^m \leq \left(\frac{b}{p} \right)^{\alpha(m-1)} C \|u\|_p^{m+\alpha p(m-1)};$$

hence (3.11) yields

$$\begin{aligned} L'(t) \geq & \left[(1-\alpha) - \frac{m-1}{m} \varepsilon k \right] H^{-\alpha}(t) H'(t) \\ & + \varepsilon \left(\frac{P}{2} + 1 \right) \int_{\Omega} u_t^2(x, t) dx + \varepsilon \left(\frac{P}{2} - 1 \right) \int_{\Omega} (\Delta u(x, t))^2 dx \\ & + \varepsilon \left[pH(t) - \frac{k^{1-m}}{m} a \left(\frac{b}{p} \right)^{\alpha(m-1)} C \|u\|_p^{m+\alpha p(m-1)} \right]. \end{aligned} \quad (3.12)$$

We then use Corollary 3.2 and relation (3.8), for $s = m + \alpha p(m - 1) \leq p$, to deduce from (3.12),

$$\begin{aligned}
 L'(t) \geq & \left[(1 - \alpha) - \frac{m - 1}{m} \varepsilon k \right] H^{-\alpha}(t) H'(t) \\
 & + \varepsilon \left(\frac{p}{2} + 1 \right) \int_{\Omega} u_t^2(x, t) dx + \varepsilon \left(\frac{p}{2} - 1 \right) \int_{\Omega} |\nabla u|^2(x, t) dx \\
 & + \varepsilon [pH(t) - C_1 k^{1-m} \{H(t) + \|u_t\|_2^2 + \|u\|_p^p\}], \tag{3.13}
 \end{aligned}$$

where $C_1 = a(b/p)^{\alpha(m-1)}C/m$. By noting that

$$H(t) = \frac{b}{p} \|u\|_p^p - \frac{1}{2} \|u_t\|_2^2 - \frac{1}{2} \|\Delta u\|_2^2$$

and writing $p = (p + 2)/2 + (p - 2)/2$, the estimate (3.13) gives

$$\begin{aligned}
 L'(t) \geq & \left[(1 - \alpha) - \frac{m - 1}{m} \varepsilon k \right] H^{-\alpha}(t) H'(t) + \varepsilon \frac{p - 2}{4} \|\Delta u\|_2^2 \\
 & + \varepsilon \left[\left(\frac{p + 2}{2} - C_1 k^{1-m} \right) H(t) + \left(\frac{p - 2}{2p} b - C_1 k^{1-m} \right) \|u\|_p^p \right. \\
 & \left. + \left(\frac{p + 6}{4} - C_1 k^{1-m} \right) \|u_t\|_2^2 \right]. \tag{3.14}
 \end{aligned}$$

At this point, we choose k large enough so that the coefficients of $H(t)$, $\|u_t\|_2^2$, and $\|u\|_p^p$ in (3.14) are strictly positive; hence we get

$$\begin{aligned}
 L'(t) \geq & \left[(1 - \alpha) - \frac{m - 1}{m} \varepsilon k \right] H^{-\alpha}(t) H'(t) \\
 & + \varepsilon \gamma [H(t) + \|u_t\|_2^2 + \|u\|_p^p], \tag{3.15}
 \end{aligned}$$

where $\gamma > 0$ is the minimum of these coefficients. Once k is fixed (hence γ), we pick ε small enough so that $(1 - \alpha) - \varepsilon k(m - 1)/m \geq 0$ and

$$L(0) = H^{1-\alpha}(0) + \varepsilon \int_{\Omega} u_0 u_1(x) dx > 0.$$

Therefore (3.15) takes the form

$$L'(t) \geq \gamma \varepsilon [H(t) + \|u_t\|_2^2 + \|u\|_p^p]. \tag{3.16}$$

Consequently we have

$$L(t) \geq L(0) > 0, \quad \forall t \geq 0.$$

Next we estimate the second term in (3.7) as follows:

$$\left| \int_{\Omega} uu_t(x, t) dx \right| \leq \|u\|_2 \|u_t\|_2 \leq C \|u\|_p \|u_t\|_2.$$

So we have

$$\left| \int_{\Omega} uu_t(x, t) dx \right|^{1/(1-\alpha)} \leq C \|u\|_p^{1/(1-\alpha)} \|u_t\|_2^{1/(1-\alpha)}.$$

Again Young's inequality gives

$$\left| \int_{\Omega} uu_t(x, t) dx \right|^{1/(1-\alpha)} \leq C \left[\|u\|_p^{\mu/(1-\alpha)} + \|u_t\|_2^{\theta/(1-\alpha)} \right], \quad (3.17)$$

for $1/\mu + 1/\theta = 1$. We take $\theta = 2(1 - \alpha)$ to get $\mu/(1 - \alpha) = 2/(1 - 2\alpha) \leq p$ by condition (3.8). Therefore (3.17) becomes

$$\left| \int_{\Omega} uu_t(x, t) dx \right|^{1/(1-\alpha)} \leq C [\|u\|_p^s + \|u_t\|_2^2],$$

where $s = 2/(1 - 2\alpha) \leq p$. By using Corollary 3.2 we obtain

$$\left| \int_{\Omega} uu_t(x, t) dx \right|^{1/(1-\alpha)} \leq C [H(t) + \|u\|_p^s + \|u_t\|_2^2], \quad \forall t \geq 0. \quad (3.18)$$

Consequently we have

$$\begin{aligned} L^{1/(1-\alpha)}(t) &= \left(H^{1-\alpha}(t) + \varepsilon \int_{\Omega} uu_t(x, t) dx \right)^{1/(1-\alpha)} \\ &\leq 2^{1/(1-\alpha)} \left(H(t) + \left| \int_{\Omega} uu_t(x, t) dx \right|^{1/(1-\alpha)} \right) \\ &\leq C(H(t) + \|u\|_p^s + \|u_t\|_2^2). \end{aligned} \quad (3.19)$$

We then combine (3.16) and (3.19), to arrive at

$$L'(t) \geq \Gamma L^{1/(1-\alpha)}(t), \quad (3.20)$$

where Γ is a constant depending on C , γ , and ε only (and hence is independent of the solution u). A simple integration of (3.20) over $(0, t)$ then yields

$$L^{\alpha/(1-\alpha)}(t) \geq \frac{1}{L^{-\alpha/(1-\alpha)}(0) - \Gamma t \alpha / (1 - \alpha)}.$$

Therefore $L(t)$ blows up in a time

$$T^* \leq \frac{1 - \alpha}{\Gamma \alpha [L(0)]^{\alpha/(1-\alpha)}}. \quad (3.21)$$

Remark 2.1. By following the steps of the proof of Theorem 3.3 closely, one can easily see that the blow-up result holds even for $1 < m < p$. Therefore this method is a unified one for both linear and nonlinear damping cases.

Remark 2.2. The estimate (3.21) shows that $L(0)$ is larger when the blow-up takes place more quickly.

4. GLOBAL EXISTENCE

In this section, we show that the solution (2.5) is global if $m \geq p$.

THEOREM 4.1. *Assume that (2.2) and (2.4) hold such that $m \geq p$. Then for any ϕ in $H_0^2(\Omega)$ and φ in $L^2(\Omega)$, the problem (1.2) has a unique weak solution $u \in \mathbf{Y}$, for any $T > 0$.*

Proof. Similar to [3], we define the functional

$$F(t) := \frac{1}{2} \int_{\Omega} [u_t^2 + (\Delta u)^2](x, t) dx + \frac{b}{p} \int_{\Omega} |u(x, t)|^p dx.$$

By taking a derivative and using Eq. (1.2), we obtain

$$F'(t) = -a \|u_t\|_m^m + 2b \int_{\Omega} u_t u |u(x, t)|^{p-2} dx.$$

By using Young's inequality, we get

$$F'(t) \leq -a \|u_t\|_m^m + \delta \|u_t\|_p^p + C_{\delta} \|u\|_p^p.$$

By noting that $m \geq p$, we easily see that

$$F'(t) \leq -a \|u_t\|_m^m + C\delta \|u_t\|_m^p + C_{\delta} \|u\|_p^p,$$

where C is a constant depending on Ω only and C_{δ} is a constant depending on δ . At this point we distinguish two cases: either $\|u_t\|_m > 1$, so we choose δ small enough so that $-a \|u_t\|_m^m + C\delta \|u_t\|_m^p \leq 0$, and hence $F'(t) \leq C_{\delta} \|u\|_p^p$. Or $\|u_t\|_m \leq 1$; in this case we have $F'(t) \leq C\delta + C_{\delta} \|u\|_p^p$. Therefore, in either case, we have

$$F'(t) \leq c_1 + cF(t).$$

A simple integration then yields

$$F(t) \leq (F(0) + c_1/c)e^{ct}.$$

The last estimate, together with the continuation principle, completes the proof.

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