

On Poisson Semigroup Generated by the Generalized B-Translation

Simten Uyhan

Akdeniz University, Department of Mathematics,
 Antalya, TURKEY

Abstract

The Poisson semigroup associated with the singular differential operator
 $\Delta_B = \sum_{k=1}^n \left(\frac{\partial^2}{\partial x_k^2} + \frac{2v_k}{x_k} \cdot \frac{\partial}{\partial x_k} \right)$ is introduced and some properties are studied.

1 Auxiliary definitions, notations and results

Suppose that \mathbb{R}^n is the n -dimensional Euclidean space, $x = (x_1, \dots, x_n)$, $\xi = (\xi_1, \dots, \xi_n)$ are vectors in \mathbb{R}^n , $(x \cdot \xi) = x_1 \xi_1 + \dots + x_n \xi_n$, $|x| = (x \cdot x)^{1/2}$ and $\mathbb{R}_+^n = \{x = (x_1, \dots, x_n); x_1 > 0, \dots, x_n > 0\}$.

We denote by $\Delta_B \equiv \Delta_B(x)$, $v = (v_1, \dots, v_n)$ the singular differential operator

$$\Delta_B = \sum_{k=1}^n \left(\frac{\partial^2}{\partial x_k^2} + \frac{2v_k}{x_k} \cdot \frac{\partial}{\partial x_k} \right) \quad (v_1 > 0, \dots, v_n > 0). \quad (1.1)$$

Let $L_{p,v} \equiv L_p(\mathbb{R}_+^n, x^{2v} dx)$, $1 \leq p < \infty$, be the space of measurable functions on \mathbb{R}_+^n with the norm

$$\|f\|_{p,v} = \left(\int_{\mathbb{R}_+^n} |f(x)|^p x^{2v} dx \right)^{1/p}; \quad x^{2v} = x_1^{2v_1} \dots x_n^{2v_n}; \quad dx = dx_1 \dots dx_n. \quad (1.2)$$

For $x \in \mathbb{R}_+^n$, $y \in \mathbb{R}_+^n$, the generalized B-translation of $f: \mathbb{R}_+^n \rightarrow C$ is defined by

$$\begin{aligned} T^y f(x) &= \pi^{-n/2} \prod_{k=1}^n \Gamma\left(v_k + \frac{1}{2}\right) \Gamma^{-1}(v_k) \int_0^\infty \dots \int_0^\infty \prod_{k=1}^n \sin^{2v_k-1} \alpha_k \quad (1.3) \\ &\times f\left(\sqrt{x_1^2 - 2x_1 y_1 \cos \alpha_1 + y_1^2}, \dots, \sqrt{x_n^2 - 2x_n y_n \cos \alpha_n + y_n^2}\right) d\alpha_1 \dots d\alpha_n. \end{aligned}$$

For the relevant one-dimensional generalized Bessel translation operator

$$S^\rho g(r) = \frac{\Gamma(v + \frac{1}{2})}{\Gamma(v)\Gamma(\frac{1}{2})} \int_0^\pi g(\sqrt{r^2 - 2r\rho \cos \alpha + \rho^2}) \sin^{2v-1} \alpha \, d\alpha$$

the following relations are known [3]:

$$\begin{aligned} S^\rho g(r) &= S^r g(\rho), & S^\rho S^\tau g(r) &= S^\tau S^\rho g(r), \\ S^\rho g(r) &= S^{-\rho} g(r), & S^0 g(r) &= g(r), \\ \int_0^\infty f(r) S^r g(\rho) r^{2v} \, dr &= \int_0^\infty S^r f(\rho) g(r) r^{2v} \, dr. \end{aligned}$$

Let $f \in L_{p,v}$, $1 \leq p < \infty$. Then for all $x \in \mathbb{R}_+^n$, the function $T^x f$ belongs to $L_{p,v}$ (see [4]) and

$$\|T^x f\|_{p,v} \leq \|f\|_{p,v}. \tag{1.4}$$

The generalized B-translation operator T^y generates the corresponding B-convolution

$$(f \otimes g)(x) = \int_{\mathbb{R}_+^n} T^y f(x) g(y) y^{2v} \, dy. \tag{1.5}$$

By using (1.4) and the Riesz-Thorin interpolation theorem it is not difficult to prove the corresponding Young inequality

$$\|f \otimes g\|_{r,v} \leq \|f\|_{p,v} \cdot \|g\|_{q,v}, \quad 1 \leq p, q, r \leq \infty, \quad \frac{1}{p} + \frac{1}{q} = \frac{1}{r} + 1. \tag{1.6}$$

The Fourier-Bessel transform and its inverse are defined by

$$(F_\nu \varphi)(z) = \int_{\mathbb{R}_+^n} \varphi(x) \left(\prod_{k=1}^n j_{\nu_k - \frac{1}{2}}(x_k z_k) \right) x^{2\nu} \, dx, \tag{1.7}$$

$$(F_\nu^{-1} \varphi)(x) = c_\nu(n) (F_\nu \varphi)(-x), \quad c_\nu(n) = \left(\prod_{k=1}^n 2^{2\nu_k} \Gamma^2\left(\nu_k + \frac{1}{2}\right) \right)^{-1}, \tag{1.8}$$

where $j_p(t)$ ($t > 0$, $p > -\frac{1}{2}$) is connected with the Bessel function of the first kind $J_p(t)$ as follows [3]:

$$j_p(t) = 2^p \Gamma(p + 1) \frac{J_p(t)}{t^p}.$$

It is known that for $f \in L_{p,v}$ ($p = 1$ or $p = 2$)

$$F_\nu(f \otimes g) = F_\nu(f) F_\nu(g). \tag{1.9}$$

2 A Poisson semigroup associated with the generalized B-translation and its properties

The Poisson semigroup associated with Δ_B is an integral operator of convolution type generated by the generalized B-translation. The kernel of this operator is defined as the Fourier-Bessel transform of the function $\exp(-\alpha|y|)$ ($y \in \mathbb{R}_+^n$, $\alpha > 0$).

Let us calculate $F_\nu(\exp(-\alpha|y|))$ by using the following formulas [5] and [1], respectively,

$$e^{-\beta} = \frac{1}{\sqrt{\pi}} \int_0^\infty \frac{e^{-t}}{\sqrt{t}} e^{-\beta^2/4t} dt,$$

$$F_\nu \left(e^{-\alpha|y|} \right) (t) = 2^{-n} \Gamma \left(\nu_1 + \frac{1}{2} \right) \dots \Gamma \left(\nu_n + \frac{1}{2} \right) \alpha^{-\frac{2\nu_1 + \dots + 2\nu_n + n}{2}} e^{-\frac{t}{4\alpha}}.$$

By Fubini's theorem we have

$$\begin{aligned} F_\nu \left(e^{-|y|} \right) (x) &= \int_{\mathbb{R}_+^n} e^{-|y|} \left(\prod_{k=1}^n j_{\nu_k - \frac{1}{2}}(x_k y_k) y_k^{2\nu_k} \right) dy \\ &= \int_{\mathbb{R}_+^n} \left(\frac{1}{\sqrt{\pi}} \int_0^\infty \frac{e^{-t}}{\sqrt{t}} e^{-|y|^2/4t} dt \right) \left(\prod_{k=1}^n j_{\nu_k - \frac{1}{2}}(x_k y_k) y_k^{2\nu_k} \right) dy \\ &= \frac{1}{\sqrt{\pi}} \int_0^\infty \frac{e^{-t}}{\sqrt{t}} \int_{\mathbb{R}_+^n} e^{-|y|^2/4t} \prod_{k=1}^n j_{\nu_k - \frac{1}{2}}(x_k y_k) y_k^{2\nu_k} dy dt \\ &= \frac{1}{\sqrt{\pi}} \int_0^\infty \frac{e^{-t}}{\sqrt{t}} \prod_{k=1}^n \int_0^\infty e^{-y_k^2/4t} j_{\nu_k - \frac{1}{2}}(x_k y_k) y_k^{2\nu_k} dy_k dt \\ &= \frac{1}{\sqrt{\pi}} \int_0^\infty \frac{e^{-t}}{\sqrt{t}} \prod_{k=1}^n \frac{1}{2} \left(\frac{1}{4t} \right)^{-\nu_k - \frac{1}{2}} \Gamma \left(\nu_k + \frac{1}{2} \right) e^{-(x_k^2/4)\frac{1}{4t}} dt \\ &= \frac{1}{\sqrt{\pi}} \int_0^\infty \frac{e^{-t}}{\sqrt{t}} \left[\frac{1}{2^n} (4t)^{\nu_1 + \frac{1}{2}} \dots (4t)^{\nu_n + \frac{1}{2}} \Gamma(\nu_1 + \frac{1}{2}) \dots \Gamma(\nu_n + \frac{1}{2}) e^{-tx_1^2} \dots e^{-tx_n^2} \right] dt \\ &= \frac{1}{\sqrt{\pi}} \frac{1}{2^n} \cdot 2^{2\nu_1 + \dots + 2\nu_n} 2^n \Gamma \left(\nu_1 + \frac{1}{2} \right) \dots \Gamma \left(\nu_n + \frac{1}{2} \right) \int_0^\infty \frac{e^{-t}}{\sqrt{t}} t^{\nu_1 + \dots + \nu_n + \frac{n}{2}} e^{-t|x|^2} dt \end{aligned}$$

$$\begin{aligned}
&= \frac{1}{\sqrt{\pi}} 2^{2v_1+\dots+2v_n} \Gamma\left(v_1 + \frac{1}{2}\right) \dots \Gamma\left(v_n + \frac{1}{2}\right) \int_0^\infty t^{\frac{n-1}{2}+v_1+\dots+v_n} e^{-t(1+|x|^2)} dt \\
&= \frac{1}{\sqrt{\pi}} 2^{v_1+\dots+v_n} \left(\sqrt{c_v(n)}\right)^{-1} \int_0^\infty t^{\frac{n+1}{2}+v_1+\dots+v_n-1} e^{-t(1+|x|^2)} dt,
\end{aligned}$$

where $c_v(n)$ is defined in (1.8).

From the last equality we get

$$\begin{aligned}
F_v\left(e^{-|y|}\right)(x) &= \frac{1}{\sqrt{\pi}} 2^{v_1+\dots+v_n} \left(\sqrt{c_v(n)}\right)^{-1} \Gamma\left(\frac{n+1}{2} + v_1 + \dots + v_n\right) \\
&\quad \times \left(1 + |x|^2\right)^{-\left(\frac{n+1}{2}+v_1+\dots+v_n\right)},
\end{aligned}$$

by using the definition of the Gamma function.

Finally, using the equality

$$F_v(f(\lambda y))(x) = \lambda^{-(n+2v_1+\dots+2v_n)} F_v(f(y))\left(\frac{x}{\lambda}\right) \quad (\lambda > 0),$$

we have

$$\begin{aligned}
F_v\left(e^{-\alpha|y|}\right)(x) &= \frac{1}{\sqrt{\pi}} 2^{v_1+\dots+v_n} \left(\sqrt{c_v(n)}\right)^{-1} \Gamma\left(\frac{n+1}{2} + v_1 + \dots + v_n\right) \\
&\quad \times \frac{\alpha}{\left(|x|^2 + \alpha^2\right)^{\frac{n+1}{2}+v_1+\dots+v_n}}.
\end{aligned}$$

In view of the last equality we define the Poisson kernel as

$$\begin{aligned}
P_v(x; \alpha) &= \sqrt{c_v(n)} \frac{1}{\sqrt{\pi}} 2^{v_1+\dots+v_n} \Gamma\left(\frac{n+1}{2} + v_1 + \dots + v_n\right) \\
&\quad \times \frac{\alpha}{\left(|x|^2 + \alpha^2\right)^{\frac{n+1}{2}+v_1+\dots+v_n}}. \tag{2.1}
\end{aligned}$$

It is not hard to verify the following properties of $P_v(x; \alpha)$:

- 1) $F_v(P_v(x; \alpha))(x) = e^{-\alpha|x|}$;
- 2) $\|P_v(\cdot; \alpha)\|_{1,v} = 1$ (for all $\alpha > 0$);
- 3) $P_v(x; \alpha + \beta) = P_v(x; \alpha) \otimes P_v(x; \beta) \equiv \int_{\mathbb{R}_+^n} P_v(y; \alpha) T_x^y(P_v(x; \beta)) y^{2v} dy$,

where the symbol T_x^y denotes the translation T^y applied to the variable x .

Now, we define the Poisson integral (semigroup) generated by the generalized translation as

$$(V_\alpha f)(x) \equiv v(x; \alpha) = \int_{\mathbb{R}_+^n} f(y) T_x^y (P_v(x; \alpha)) y^{2v} dy. \tag{2.2}$$

By making use of the Fourier-Bessel transform, it is not difficult to verify that the Poisson integral $v(x; \alpha)$ is the solution of the following boundary value problem

$$\begin{cases} \left(\frac{\partial^2}{\partial \alpha^2} + \Delta_B(x) \right) v(x; \alpha) = 0, \\ v(x; \alpha) |_{\alpha=0} = f(x) \end{cases}$$

for “good” f .

It is easy to show the semigroup property, $V_\alpha V_\beta = V_{\alpha+\beta}$ ($0 < \alpha, \beta < \infty$) of $\{V_\alpha\}$, $\alpha > 0$ by using the Fourier-Bessel transform

$$F_v(V_{\alpha+\beta}f) = e^{-(\alpha+\beta)|y|} F_v f = e^{-\alpha|y|} \left(e^{-\beta|y|} F_v f \right) = F_v(V_\alpha V_\beta f).$$

Now, let us investigated the approximation and other properties of the Poisson semigroup $V_\alpha f$ ($\alpha > 0$) associated with Δ_B . For this, first we define the Hardy-Littlewood maximal function $M_B f$ generated by the generalized translation. The maximal function $M_B f$ of $f \in L_{p,v}$, $1 \leq p \leq \infty$, is defined as

$$M_B f(x) = \sup_{r>0} \frac{1}{|E_+(0,r)|} \int_{E_+(0,r)} T^y |f(x)| y^{2v} dy,$$

where $E_+(x,r) = \{y : y \in \mathbb{R}_+^n, |x - y| < r\}$,

$$|E_+(0,r)| = \int_{E_+(0,r)} y^{2v} dy = w(n,v) r^{n+2v_1+\dots+2v_n}, \quad w(n,v) = \int_{E_+(0,1)} x^{2v} dx.$$

It is known (see [2]) that the maximal operator M_B is of weak type (1,1) and is bounded on $L_{p,v}$, $1 < p \leq \infty$.

We can prove the following lemma for the maximal function $M_B f$ by using some ideas in E. Stein and G. Weiss’s monograph [5].

Lemma 1 *Let $\varphi \in L_{1,v}$ be a radial and $\psi(r) = \varphi(x) |_{|x|=r}$ ($0 < r < \infty$) be a non-negative and decreasing function on $[0, \infty)$. Then for every $f \in L_{p,v}$ ($1 \leq p \leq \infty$)*

we have

$$\sup_{\varepsilon>0} |(f \otimes \varphi_\varepsilon)(x)| \leq \|\varphi\|_{1,v} M_B f(x), \tag{2.3}$$

where $\varphi_\varepsilon(x) = \varepsilon^{-n-(2v_1+\dots+2v_n)} \varphi\left(\frac{x}{\varepsilon}\right)$ ($\varepsilon > 0$).

Proof. For the sake of natural simplicity, we assume $f \geq 0$.

Step I. Let the function φ be defined by

$$\varphi(x) = \begin{cases} \frac{1}{w(v,n)} & , \quad x \in E_+(0,1), \\ 0 & , \quad x \in \mathbb{R}_+^n \setminus E_+(0,1). \end{cases}$$

We have $\|\varphi\|_{1,v} = 1$. Putting $\varphi_\varepsilon(x) = \varepsilon^{-n-(2v_1+\dots+2v_n)}\varphi\left(\frac{x}{\varepsilon}\right)$, we get $\|\varphi_\varepsilon\|_{1,v} = \|\varphi\|_{1,v} = 1$ for all $\varepsilon > 0$. Then

$$M_B f(x) = \sup_{\varepsilon > 0} |(f \otimes \varphi_\varepsilon)(x)| \quad \text{for all } f \in L_{p,v} \quad (f \geq 0).$$

Step II. Let $\varphi(x) = \sum_{k=1}^m c_k \chi_k(x)$ ($c_k \geq 0$, $k = 1, \dots, m$), where $\chi_k(x)$ is the characteristic function of the sphere $E_+(0, r_k)$. Putting $\varphi_\varepsilon(x) = \varepsilon^{-n-(2v_1+\dots+2v_n)}\varphi\left(\frac{x}{\varepsilon}\right)$, we get

$$\begin{aligned} (f \otimes \varphi_\varepsilon)(x) &= \sum_{k=1}^m c_k \varepsilon^{-n-(2v_1+\dots+2v_n)} \int_{E_+(0, \varepsilon r_k)} T^y f(x) y^{2v} dy \\ &= \sum_{k=1}^m c_k w(n, v) r_k^{n+(2v_1+\dots+2v_n)} \frac{1}{w(n, v) (\varepsilon r_k)^{n+(2v_1+\dots+2v_n)}} \int_{E_+(0, \varepsilon r_k)} T^y f(x) y^{2v} dy \\ &\leq M_B f(x) \sum_{k=1}^m c_k w(n, v) r_k^{n+(2v_1+\dots+2v_n)} \\ &= M_B f(x) \sum_{k=1}^m c_k \int_{E_+(0, r_k)} x^{2v} dx \\ &= M_B f(x) \sum_{k=1}^m c_k \int_{\mathbb{R}_+^n} \chi_k(x) x^{2v} dx \\ &= M_B f(x) \int_{\mathbb{R}_+^n} \left(\sum_{k=1}^m c_k \chi_k(x) \right) x^{2v} dx \\ &= M_B f(x) \|\varphi\|_{1,v}. \end{aligned}$$

Thus,

$$\sup_{\varepsilon > 0} |(f \otimes \varphi_\varepsilon)(x)| \leq \|\varphi\|_{1,v} M_B f(x)$$

for every nonnegative simple function φ .

Step III. Since $\psi(r)$ is nonnegative decreasing on $[0, \infty)$ and the function $\varphi \in L_{1,v}$ is of the form $\varphi(x) = \psi(|x|)$, then we have $\psi(r) \rightarrow 0$ as $r \rightarrow \infty$. Thus, it is possible to approximate the nonnegative function $\varphi(x) = \psi(|x|)$ from below by an increasing sequence of simple functions of the type $\varphi_m(x) = \sum_{k=1}^m c_k^m \chi_k(x)$. We have proved above the inequality (2.3) for the simple functions φ_m . Now, taking the limit as $m \rightarrow \infty$, one concludes the proof. ■

The following theorem states the main result of this work which gives some properties of the Poisson integral generated by the Laplace-Bessel differential operator Δ_B .

Theorem 2 *Let $V_\alpha f$ ($\alpha > 0$) be the Poisson semigroup for a function f . If $f \in L_{p,v}$, $1 \leq p \leq \infty$, then:*

- a) $\|V_\alpha f\|_{p,v} \leq \|f\|_{p,v}$;
- b) $\sup_{\alpha > 0} |(V_\alpha f)(x)| \leq M_B f(x)$;
- c) $V_\alpha V_\beta f = V_{\alpha+\beta} f$ ($\alpha > 0, \beta > 0$) ;
- d) $\operatorname{ess\,sup}_{x \in \mathbb{R}_+^n} |(V_\alpha f)(x)| \leq C \alpha^{-\frac{n+2v_1+\dots+2v_n}{p}} \|f\|_{p,v}$;
- e) $(L_{p,v}) \lim_{\alpha \rightarrow 0^+} V_\beta f = f$ ($1 \leq p < \infty$),

where $(L_{p,v}) \lim$ denotes the limit in the norm $L_{p,v}$ and pointwise for almost all $x \in \mathbb{R}^n$.

Proof.

a) By using Young's inequality (1.6) and the equality $\|P_v(\cdot; \alpha)\|_{1,v} = 1$ for all $\alpha > 0$, we have $\|V_\alpha f\|_{p,v} = \|f \otimes P_v(\cdot; \alpha)\|_{p,v} \leq \|f\|_{p,v} \|P_v(\cdot; \alpha)\|_{1,v}$.

b) The proof of this result follows directly from (2.3) by taking $\varphi_\varepsilon(x) = P_v(x; \varepsilon) \equiv \varepsilon^{-n-(2v_1+\dots+2v_n)} P_v\left(\frac{x}{\varepsilon}; 1\right)$.

c) is proved above.

d) is obtained by substituting ∞ for r and $P_v(x; \alpha)$ for the function g in Young's inequality (1.6) and using the homogeneity property of $P_v(\cdot; \alpha)$.

e) By using the equality $\|P_\nu(\cdot; \alpha)\|_{1,\nu} = 1$ we have

$$\|V_\alpha f - f\|_{p,\nu} = \left\| \int_{\mathbb{R}_+^n} P_\nu(x; \alpha) [T^x f(y) - f(y)] x^{2\nu} dx \right\|_{p,\nu}.$$

By setting αx instead of x and applying the generalized Minkowski's inequality, we have

$$\begin{aligned} \|V_\alpha f - f\|_{p,\nu} &= \left\| \int_{\mathbb{R}_+^n} P_\nu(x; 1) [T^{\alpha x} f(y) - f(y)] x^{2\nu} dx \right\|_{p,\nu} \\ &\leq \int_{\mathbb{R}_+^n} P_\nu(x; 1) \| [T^{\alpha x} f(y) - f(y)] \|_{p,\nu} x^{2\nu} dx \end{aligned}$$

by using (1.4) and the property

$$\lim_{t \rightarrow 0} \|T^t f(\cdot) - f(\cdot)\|_{p,\nu} = 0 \quad (\text{see [4]})$$

of T^t and, finally, applying the Lebesgue dominated convergence theorem we get

$$\lim_{\alpha \rightarrow 0} \|V_\alpha f - f\|_{p,\nu} = 0. \quad \blacksquare$$

References

- [1] ALIEV I. A. AND BAYRAKÇI S., *On inversion of B-elliptic potentials by the method of Balakrishnan-Rubin*, Fractional Calculus and Applied Analysis, **1** (1998), No. 4, 365–384.
- [2] GULIEV V. S., *On maximal function and fractional integral, associated with the Bessel differential operator*, Mathematical Inequalities and Applications, **6** (2003), No. 2, 137–330.
- [3] LEVITAN B. M., *Bessel functions expansions in series and Fourier integrals*, Uspekhi Mat. Nauk, **6** (1952), No. 2, 102–143.
- [4] LÖFSTRÖM J. AND PEETRE J., *Approximation theorems connected with generalized translation*, Math. Ann., **181** (1969), 255–268.
- [5] STEIN E. M. AND WEISS G., *Introduction to Fourier Analysis on Euclidean Spaces*, Princeton Univ. Press, Princeton, N.J., 1971.
- [6] STEIN E. M., *Singular Integrals and Differentiability Properties of Functions*, Princeton Univ. Press, Princeton, N.J., 1970.