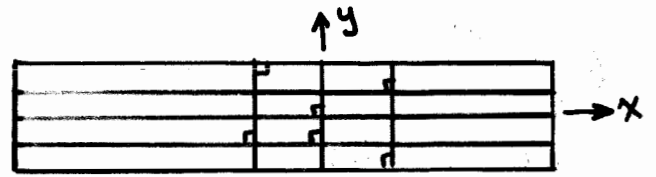


Pure Bending of Prismatic Beams

* The basic 3 ingredients of mechanics are employed here.

* Kinematic Consideration and assumptions:

- ① Transverse cross sections \perp the centroidal line of the beam before bending remain plane and \perp to the centroidal line after deformation (loading)



← Bernoulli hypothesis

- ② Cross sections are not assumed rigid.
 \Rightarrow Lengths of lines differ after deformation.

$$\sigma_y = \sigma_z = 0$$

$$\Rightarrow \epsilon_y = \epsilon_z = -\nu \epsilon_x$$

The assumptions above are geometric, but limited to small deformation.

① Compat.

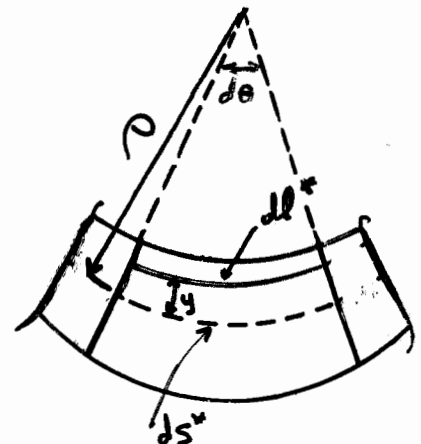
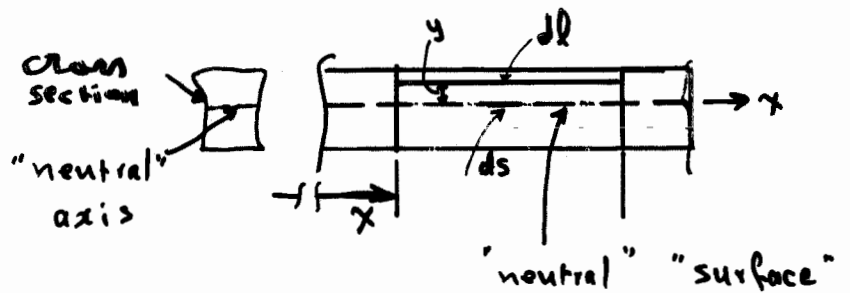
From the figure:

$$\begin{aligned} \epsilon_x &= \frac{\Delta l}{l} \\ &= \frac{\Delta l^* - \Delta l}{\Delta l} \end{aligned}$$

Note that

$$\begin{aligned} ds &= \Delta l \\ ds^* &= \Delta l^* \\ \Delta l^* &\neq \Delta l \end{aligned}$$

$$\begin{aligned} \Rightarrow \epsilon_x &= \frac{\Delta l^* - ds}{ds} \\ &= \frac{\Delta l^*}{ds} - 1 \end{aligned}$$



From similar triangles (assuming "small" $d\theta \Rightarrow ds^*$ and $dl^* \approx$ straight lines)

$$\Rightarrow \frac{\Delta l^*}{\rho - y} = \frac{ds^*}{\rho} = \frac{ds}{\rho}$$

$$\Rightarrow \frac{d l^0}{d s} = 1 - \frac{y}{\rho} \Rightarrow \frac{d l^0}{d s} - 1 = -\frac{y}{\rho} \Rightarrow \underline{\underline{\epsilon_x = -\frac{y}{\rho}}}$$

ρ = radius of curvature of the deformed centroidal line at section x .

Note that ϵ_x varies linearly through the depth of the beam.

* The assumptions above are geometric and independent of material behavior.

② * Equilibrium

Note that for elastic isotropic materials: $\sigma_y = \sigma_z = \tau_{xy} = \tau_{xz} = \tau_{yz} = 0$

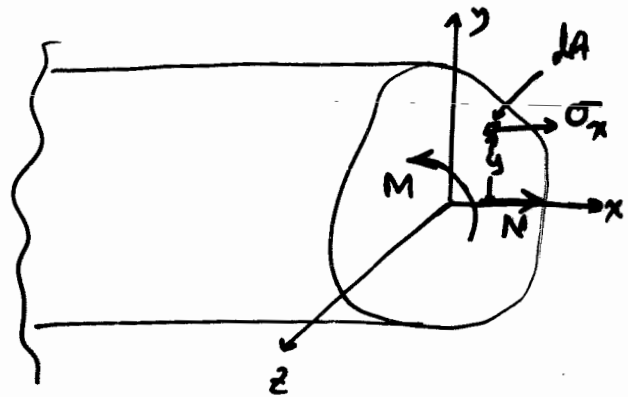
\Rightarrow Only σ_x is nonzero (pure bending) which is normal to the cross-section.

Thus, for a cross-section

$$N = \text{normal force} = \int_{\text{Area}} \sigma_x dA$$

and

$$M = - \int_A \sigma_x y dA$$



The equations above are for any material.

③ * Material Behavior

Assume linear elastic behavior \Rightarrow

$$\sigma_x = E \epsilon_x = E \left(-\frac{y}{\rho} \right) = -\frac{E y}{\rho}$$

$$M = - \int_A \sigma_x y dA = - \int_A \left(-\frac{E y}{\rho} \right) y dA$$

$$= \int_A \frac{E}{\rho} y^2 dA$$

$$= \frac{E}{\rho} \int_A y^2 dA$$

Recall from Statics :

$$\int_A y^2 dA = I$$

= moment of inertia with respect to the neutral axis (N.A.)

$$\Rightarrow M = \frac{EI}{\rho} \Rightarrow \underline{\underline{\frac{1}{\rho} = \frac{M}{EI}}} \leftarrow \text{moment-curvature relationship}$$

From above, $\sigma_x = -\frac{E y}{\rho}$

$$= -E y \left(\frac{1}{\rho} \right)$$

$$= -E y \left(\frac{M}{EI} \right) \Rightarrow$$

$$\sigma_x = -\frac{M y}{I}$$

flexure formula

Note that y is measured from the N.A. \Rightarrow We need to locate the N.A.

$$N = \int_A \sigma dA = -\frac{E}{\rho} \int_A y dA = 0 \quad \left(\text{pure bending} \Rightarrow \right)$$

$$\frac{E}{\rho} \neq 0 \quad \text{in general}$$

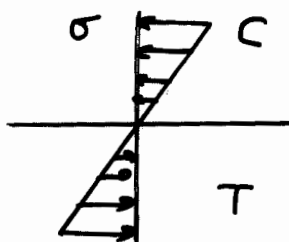
$$\Rightarrow \int_A y dA = 0$$

= first moment of area wrt the N.A.

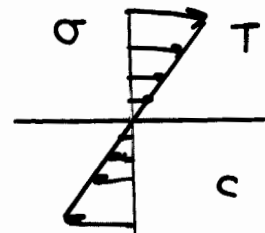
Since it is zero, this means it's the first moment of the area wrt the centroidal axis. \Rightarrow

The neutral axis coincides with the centroidal axis of the cross section for pure elastic bending.

Stress Distribution



+ M



- M

